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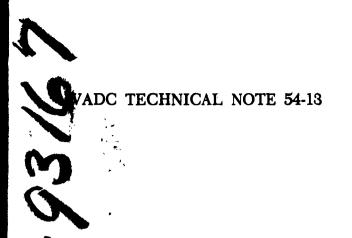
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# GENERALIZED EQUATIONS AND PROCEDURES FOR THE CALCULATION OF DETONATION PARAMETERS

Case I. Ideal Gaseous Mixtures

ROBERT G. DUNN BERNARD T. WOLFSON

AERONAUTICAL RESEARCH LABORATORY

**MARCH 1956** 

WRIGHT AIR DEVELOPMENT CENTER

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TASK No. 70164 PROJECT No. 3012

WRIGHT AIR DEVELOPMENT CENTER
AIR RESEARCH AND DEVELOPMENT COMMAND
UNITED STATES AIR FORCE
WRIGHT-PATTERSON AIR FORCE BASE, OHIO

### FOREWORD

This report was prepared by Mr. Robert G. Dunn and Mr. Bernard T. Wolfson of the Aeronautical Research Laboratory, WADC, under Task No. 70164, Project No. 3012, entitled "Ramjet Technology."

The authors wish to acknowledge the assistance given by the members of the Fluid Dynamics Research Branch, Aeronautical Research Laboratory, during the preparation of this report. In particular, gratitude is expressed to Dr. D. G. Samaras for his advice, criticism and suggestions, to Mrs. Amy Matonis and Miss Mary L. Hudson, who carried out a large portion of the calculations and assisted in the preparation of the graphs.

### ABSTRACT

The classical equations of detonation in gaseous mixtures are presented in generalized form involving only dimensionless quantities. This report sets forth a comprehensive procedure for the calculation of detonative Mach numbers and associated parameters for any combustible gaseous mixture.

The use of the method of calculation developed is illustrated by means of a sample calculation. For this purpose, calculations are presented for determination of detonation parameters for a stoichiometric mixture of hydrogen and oxygen at an initial temperature and pressure of 18°C and atmospheric pressure, respectively.

The calculation method developed is readily adaptable to computations by automatic computing machines as well as by manually-operated machines.

Graphs of basic data are presented for use in the sample calculations, and it is believed that these data will find practical use in other applications as well.

### PUBLICATION REVIEW

This report has been reviewed and is approved.

FOR THE COMMANDER:

ALDRO LINGARD Colonel, USAF

Chief, Aeronautical Research

Laboratory

Directorate of Research

alder Lingard

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### LIST OF SYMBOLS

SYMBOL	DEFINITION	UNITS
a	Velocity of sound	cm/sec
<b>a</b> 1	Velocity of sound in the initial (unburned) mixture	cm/sec
<b>a</b> 2	Velocity of sound in the final (burned) mixture	cm/sec
В	Energy release function (B = $\frac{Jh}{a^2}$ )	dimensionless
A, B, C, D, E,	Pertaining to components A, B, C, D, E, respectively, in the initial mixture	
c <sub>11</sub> , c <sub>12</sub> , c <sub>13</sub> ,	Miscellaneous constants	
C <sub>p</sub>	Point or true molal heat capacity at constant pressure	cal/gm mole
c <sub>p1</sub>	Molal heat capacity at constant pressure for the initial mixture at the initial conditions (State 1)	cal/gm mole, *K
c <sub>p2</sub>	Molal heat capacity at constant pressure for the final mixture at the final conditions (State 2)	cal/gm mole, °K
$C_{p(U)}, C_{p(V)}, C_{p(W)},$ $C_{p(X)}, C_{p(Y)}, C_{p(Z)}$	Molal heat capacities of components U, V, W, X, Y, Z, respectively in the final mixture	cal/gm mole, *K
$\overline{c}_{p}$	Mean molal heat capacity	cal/gm mole, *K
$C_{\mathbf{p}(\mathbf{i})}^{\bullet}$	Molal heat capacity at constant pres- sure for a component "i" in the ini- tial mixture at the initial conditions (State 1)	cal/gm mole, °K
C <sub>p(j)</sub>	Molal heat capacity at constant pressure for a component "j" in the final mixture at the final conditions (State 2)	cal/gm mole, °K
c <sub>V</sub>	Specific heat at constant volume	cal/gm, *K

SYMBOL	DEFINITION	UNITS
c <sub>v</sub> 1	Specific heat at constant volume for the initial mixture at the initial conditions (State 1)	cal/gm, °K
c <sub>v2</sub>	Specific heat at constant volume for the final mixture at the final conditions (State 2)	cal/gm, °K
$C_{\mathbf{v}}$	Molal heat capacity at constant volume	cal/gm mole, •K
c <sub>v1</sub>	Molal heat capacity at constant volume for the initial mixture at the initial conditions	cal/gm mole, •K
C <sub>v2</sub>	Molal heat capacity at constant volume for the final mixture at the final conditions (State 2)	cal/gm mole, •K
¯C <sub>v</sub>	Mean molal heat capacity at constant volume	cal/gm mole, •K
¯c <sub>v</sub> <sub>1,0</sub>	Mean specific heat at constant volume of initial mixture between State 1 (initial conditions) and State "0" (Reference State)	cal/gm, •K
¯c <sub>v</sub> 2, 0	Mean specific heat at constant volume of final mixture between State 2 (final conditions) and State "0" (Reference State)	cal/gm,°K
D	Detonation velocity, i.e., velocity of the detonation wave with respect to the initial mixture	cm/sec
e	Specific internal energy	cal/gm
e°	Specific internal energy of a mixture at the reference state (State "0")	cal/gm
e <sub>1</sub>	Specific internal energy of the initial mixture at the initial conditions (State 1)	cal/gm

SYMBOL	DEFINITION	UNITS
•2	Specific internal energy of the final mixture at the final conditions (State 2)	cal/gm
F,G	Symbols representing elements contained in a chemical compound	
h	Energy release per unit mass of mixture	cal/gm
H	Enthalpy	cal/gm mole
$H_i^t$	Enthalpy of compound "i" at the initial conditions (i.e., at t <sub>1</sub> ) relative	cal/gm mole
	to its standard state at temperature tL	
H <sup>!</sup> j	Enthalpy of compound "j" at the initial conditions (i.e., at t <sub>1</sub> ) relative to	cal/gm mole
	it standard state at temperature t	
△ H <sub>f</sub>	Heat of formation (The heat of formation of a compound represents the heat absorbed or evolved when one mole of the compound is formed from the elements in a reaction beginning and ending at some given temperature, with the reacting elements originally in the states of aggregation which are stable at that temperature).	cal/gm mole
ΔH <sup>1</sup> <sub>f(i)</sub>	Heat of formation of any component "i" in the initial mixture at the initial conditions (i.e., at temperature t <sub>1</sub> )	cal/gm mole
$\Delta H_{f(j)}^{i}$	Heat of formation of any component "j" in the final mixture at the initial conditions (i.e., at temperature t <sub>1</sub> )	cal/gm mole
$\Delta_{H'f(A)}$ , $\Delta_{H'f(E)}$ , $\Delta_{H'f(C)}$ , $\Delta_{H'f(D)}$ , $\Delta_{H'f(E)}$	Heats of formation of components A, B, C,D, E, respectively, in the initial mixture at the initial conditions (i.e., at temperature t <sub>1</sub> )	cal/gm mole

SYMBOL	DEFINITION	UNITS
$\Delta H_{f(\mathbf{U})}^{i}$ , $\Delta H_{f(\mathbf{V})}^{i}$ , $\Delta H_{f(\mathbf{W})}^{i}$ , $\Delta H_{f(\mathbf{Z})}^{i}$	Heats of formation of components $U, V, W, X, Y, Z$ , respectively, in the final mixture at the initial conditions (i.e., at temperature $t_1$ )	cal/gm mole
ΔH <sub>f<sub>L</sub>(j)</sub>	Heat of formation of any component "j" in the final mixture at some temperature $t_L$ , as found in the literature	cal/gm mole
ΔH <sub>f<sub>L</sub>(i)</sub>	Heat of formation of any component "i" in the initial mixture at some temperature $t_L$ , as found in the literature	cal/gm mole
ΔH <sub>R</sub>	Heat of reaction, energy release per mole of initial mixture	cal/gm mole
$\sum H_{e(i)}^{i}$	Sum of the enthalpies (at t <sub>1</sub> ) of the elements involved in the formation of one gram mole of compound "i", relative to the standard states of the elements at t <sub>L</sub>	cal/gm mole of "i"
$\sum H_{e(j)}^{i}$	Sum of the enthalpies (at t <sub>1</sub> ) of the elements involved in the formation of one gram mole of compound "j", relative to the standard states of the elements at t <sub>L</sub>	cal/gm mole of "j"
J	Mechanical equivalent of heat; J=4.185 x 10	dyne-cm/cal or ergs/cal
k <sub>11</sub> , k <sub>12</sub> , k <sub>13</sub> ,	Miscellaneous constants	
к <sub>р</sub>	Equilibrium constant, in terms of partial pressures, for the dissociation of a compound (Units for this constant are atmospheres raised to a power which is dependent upon the particular dissociation reaction considered).	(atmos) exp.
$K_{p(j)}, K_{p(j)}, K_{p(j)}, \dots$	Equilibrium constants in terms of partial pressures for different dissociation reactions for any component "j" in the final mixture	(atmos) exp.

SYMBOL	DEFINITION	UNITS
M	Mach number; $M = \frac{V}{a}$	dimensionless
M <sub>1</sub>	Detonative Mach number, i.e., the Mach number of the detonation wave with respect to the initial mixture $(M_1 = \frac{V_1}{a_1})$	dimensionless
M <sub>2</sub>	Mach number of the detonation wave with respect to the final mixture $(M_2 = \frac{V_2}{a_2})$	dimensionless
M	Pseudo-Mach number, i.e., a very close approximation to the detonative Mach number, $M_1$ ; used for convenience only and defined by $\overline{M}_1^2 = M_1^2 + \frac{\left(\frac{F_1}{\Gamma}\right)^2}{M_1^2}$	dimensionless
M <sub>w</sub> (A), M <sub>w</sub> (B), M <sub>w</sub> (C), M <sub>w</sub> (D), M <sub>w</sub> (E)	Molecular weights of components A, B, C, D, E, respectively, in the initial mixture	gms/gm mole
$M_{\mathbf{w}_1}$	Average molecular weight of the initial mixture	gms/gm mole
M <sub>w</sub> <sub>2</sub>	Average molecular weight of the final mixture	gms/gm mole
<sup>M</sup> w(i)	Molecular weight of any component "i" in the initial mixture	gms/gm mole
M <sub>w(j)</sub>	Molecular weight of any component "j" in the final mixture	gms/gm mole
$M_{w(U)}, M_{w(V)}, M_{w(W)},$ $M_{w(X)}, M_{w(Y)}, M_{w(Z)}$	Molecular weight of components U, V, W, X, Y, Z, respectively, in the final mixture	gms/gm mole
n	Number of moles	

SYMBOL	DEFINITION	UNITS
n <sub>1</sub>	Total number of moles in initial mixture	
<sup>n</sup> 2	Total number of moles in final mixture	
$\mathbf{n_i^t}$	Number of moles of any component "i" in the initial mixture	
<sup>n</sup> j	Number of moles of any component "j" in the final mixture	
n'A, n'B, n'C, n'D, n'E	Number of moles of components A, B, C, D, E, respectively, in the initial mixture	
<sup>n</sup> u, <sup>n</sup> v, <sup>n</sup> w, <sup>n</sup> x, <sup>n</sup> y, <sup>n</sup> z	Number of moles of components U, V, W X, Y, Z, respectively, in the final mixture	
N	Number of atoms in one molecule of a given component of a mixture (al- ways an integral value or zero)	
N <sub>i</sub> , N <sub>i</sub> ,	Number of atoms of element F, G, respectively, in one molecule of any component "i" in the initial mixture (always an integral value or zero)	
'N <sub>j</sub> , 'N <sub>j</sub> ,	Number of atoms of element F, G,, respectively, in one molecule of any component "j" in the final mixture (always an integral value or zero)	
p	Pressure	atmos
$\mathbf{p_1}$	Total pressure for the initial mixture at the initial conditions (State 1)	atmos
P <sub>2</sub>	Total pressure for the final mixture at the final conditions (State 2)	atmos

SYMBOL	DEFINITION	UNITS
$\mathbf{p_i^t}$	Partial pressure of any component "i" in the initial mixture at the initial conditions (State 1)	atmos
<sup>P</sup> j	Partial pressure of any component "j" in the final mixture at the final conditions (State 2)	atmos
Pu, Pv, Pw, Px, Py, Pz	Partial pressures of components U, V, W, X, Y, Z, respectively, in the final mixture at the final conditions (State 2)	atmos
P	Pressure (1.0133 x 10 <sup>6</sup> dynes/cm <sup>2</sup> = 1 atmo	os) dynes/cm <sup>2</sup>
P <sub>1</sub>	Total pressure of the initial mixture at the initial conditions (State 1)	dynes/cm <sup>2</sup>
P <sub>2</sub>	Total pressure of the final mixture at the final conditions (State 2)	dynes/cm <sup>2</sup>
q	Heat added to the system	cal/gm
R	Universal gas constant; R=1.98719	cal/gm mole, *K.
R <sub>1</sub>	Specific gas constant for initial mix- ture at the initial conditions (State 1)	cal/gm, °K
	$(R_1 = \frac{R}{M_{W_1}})$	
R <sub>2</sub>	Specific gas constant for final mix- ture at the final conditions (State 2)	cal/gm, °K
	$(R_2 = \frac{R}{M_{W_2}})$	
t	Temperature	°C
<sup>t</sup> L	A reference temperature pertaining to data found in the literature	°C
<sup>t</sup> 1	Temperature at the initial conditions (State 1)	•c

SYMBOL	DEFINITION	UNITS
<sup>t</sup> 2	Temperature at the final conditions (State 2)	•c
T	Absolute temperature	•K
T <sub>0</sub>	Datum or reference temperature	°K
т	Temperature at the initial conditions (State 1)	°K
т <sub>2</sub>	Temperature at the final conditions (State 2)	°K
U, V, W, X, Y, Z	Pertaining to components U, V, W, X, Y, Z, respectively, in the final mixture	
v	Specific volume, i.e., volume per unit mass	cm <sup>3</sup> /gm
v	Velocity	cm/sec
$\mathbf{v}_1$	Velocity of detonation wave with respect to the initial mixture	cm/sec
v <sub>2</sub>	Velocity of detonation wave with respect to the final mixture	cm/sec
w	Work done by the system	ergs or(dyne-cm) gm
x	Mole fraction	
x'A, x'B, x'C, x'D, x'E	Mole fractions of components A, B, C, D, E, respectively, in the initial mixture	
x¦	Mole fraction of any component "i" in the initial mixture	
* <sub>j</sub>	Mole fraction of any component "j" in the final mixture	

SYMBOL	DEFINITION	UNITS
*u',*v',*w',*x',*y',*z	Mole fractions of components U, V, W, X, Y, Z, respectively, in the final mixture	
E	Dimensionless parameter (epsilon) used for convenience only and defined by $\mathbf{E} = (\mathbf{Y}_2^2 - \mathbf{I}) \mathbf{B}$	dimensionless
Γ	Gamma Function, a dimensionless parameter used for convenience only and defined by $\Gamma = \frac{V_1 - V_1}{V_1 - V_1}$	dimensionless
θ	Symbol designating (or signifying) any parameter	
Δ	Symbol indicating difference	
8	Point or true specific heat (or molal heat capacity) ratio for a mixture $ (                                  $	dimensionless
<b>δ</b> ,	Specific heat (or molal heat capacity) ratio of the initial mixture at the initial conditions (State 1)	dimensionless
	(	
₹2	Specific heat (or molal heat capacity) ratio of the final mixture at the final conditions (State 2)	dimensionless
	$(\delta_2 = C_{p_2}/C_{v_2})$	
P	Density, i.e., mass per unit volume	gm/cm <sup>3</sup>
$\boldsymbol{\rho}_{i}$	Density of initial mixture at the initial conditions (State 1)	gm/cm <sup>3</sup>
$\rho_{2}$	Density of final mixture at the final conditions (State 2)	gm/cm <sup>3</sup>

SYMBOLS	DEFINITION	UNITS
Σ	Symbol indicating summation	
Subscripts		
1	Pertaining to the initial mixture at the initial conditions (State 1)	
2	Pertaining to the final mixture immediately behind the detonation wave at the final conditions (State 2)	
A, B, C,D, E	Pertaining to components A, B, C, D, E, respectively, in the initial mixture	
i	Pertaining to any component "i" in the initial mixture	
j	Pertaining to any component "j" in the final mixture	
L	Pertaining to data obtained from the literature	
0	Pertaining to a reference condition	
p	Pertaining to constant pressure	
v	Pertaining to constant volume	
U, V, W, X, Y, Z	Pertaining to components U, V, W, X, Y, Z, respectively, in the final mixture	
Left Superscripts		
F, G,	Pertaining to elements F, G,, respectively	
Right Superscripts		
"(*)" (Degree sign)	Pertaining to the standard state	

### L INTRODUCTION

The literature on the subject of detonation in gaseous mixtures dates back more than seventy years 16 17.18, 19.20, 21.22, 23, 24. The majority of this literature has emphasized the prevention of detonation rather than its utilization. In connection with combustion in propulsive devices, the study of detonation prevention is important because of the destructive effects which may be associated with uncontrolled detonative combustion. On the other hand, useful application of detonative type combustion in propulsive devices appears to warrant consideration.

In the normal combustion processes occurring in the combustion chambers of present-day combustion devices, the flame speeds are on the order of 1 to 2 feet per second, whereas, rates of propagation for stable detonative type combustion are as high as 3000 to 12,000 feet per second for various gaseous mixtures.

In order to be able to reasonably evaluate the possible advantages of detonative type combustion in propulsion applications, it is necessary to be able to determine theoretically the values of all detonation parameters for various gaseous mixtures considered, under specified initial conditions. Determination of these detonation parameters can be accomplished experimentally; however, the necessary equipment is highly specialized and the techniques involved are not sufficiently developed as yet to permit the obtaining of values for these parameters readily under all desired initial conditions.

Methods of calculating detonation parameters utilizing directly the classical detonation equations have been reported by Zeldovich and Ratner  $\frac{14}{}$  and Edse 15/.

In performing calculations of detonation parameters for gaseous mixtures, it was found to be highly desirable to rearrange the classical (Chapman-Jouguet) detonation equations into generalized dimensionless form. In this report, the velocity of detonation is expressed in the form of the Mach number of the detonation wave relative to the unburned gases; the heat of reaction is expressed in the form of a dimensionless energy release function. The other detonation parameters are all expressed as dimensionless ratios of the quantities before and after the detonation wave, i.e., pressure ratio, temperature ratio, density ratio and molecular weight ratio. Several other dimensionless quantities are introduced to simplify further the graphical presentation of the generalized data and to facilitate calculations.

A method is presented in this report for the calculation of detonation velocities and other detonation parameters for gaseous mixtures. The use of the method of calculation developed is illustrated by means of a sample calculation. For this purpose, calculations are presented for determination of the detonation parameters for a stoich-iometric mixture of hydrogen and oxygen at an initial temperature and pressure of 18°C and atmospheric pressure, respectively.

MOTE - This Technical Note was released by the author for publication in December 1954.

### L INTRODUCTION

The literature on the subject of detonation in gaseous mixtures dates back more than seventy years 16 17.18.19.20.21.22.23.24. The majority of this literature has emphasized the prevention of detonation rather than its utilization. In connection with combustion in propulsive devices, the study of detonation prevention is important because of the destructive effects which may be associated with uncontrolled detonative combustion. On the other hand, useful application of detonative type combustion in propulsive devices appears to warrant consideration.

In the normal combustion processes occuring in the combustion chambers of present-day combustion devices, the flame speeds are on the order of 1 to 2 feet per second, whereas, rates of propagation for stable detonative type combustion are as high as 3000 to 12,000 feet per second for various gaseous mixtures.

In order to be able to reasonably evaluate the possible advantages of detonative type combustion in propulsion applications, it is necessary to be able to determine theoretically the values of all detonation parameters for various gaseous mixtures considered, under specified initial conditions. Determination of these detonation parameters can be accomplished experimentally; however, the necessary equipment is highly specialized and the techniques involved are not sufficiently developed as yet to permit the obtaining of values for these parameters readily under all desired initial conditions.

Methods of calculating detonation parameters utilizing directly the classical detonation equations have been reported by Zeldovich and Ratner  $\frac{14}{}$  and Edse  $\frac{15}{}$ .

In performing calculations of detonation parameters for gaseous mixtures, it was found to be highly desirable to rearrange the classical (Chapman-Jouguet) detonation equations into generalized dimensionless form. In this report, the velocity of detonation is expressed in the form of the Mach number of the detonation wave relative to the unburned gases; the heat of reaction is expressed in the form of a dimensionless energy release function. The other detonation parameters are all expressed as dimensionless ratios of the quantities before and after the detonation wave, i.e., pressure ratio, temperature ratio, density ratio and molecular weight ratio. Several other dimensionless quantities are introduced to simplify further the graphical presentation of the generalized data and to facilitate calculations.

A method is presented in this report for the calculation of detonation velocities and other detonation parameters for gaseous mixtures. The use of the method of calculation developed is illustrated by means of a sample calculation. For this purpose, calculations are presented for determination of the detonation parameters for a stoich-iometric mixture of hydrogen and oxygen at an initial temperature and pressure of 18°C and atmospheric pressure, respectively.

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### IL DERIVATION OF THE GENERALIZED EQUATIONS IN NON-DIMENSIONAL FORM

Some of the assumptions utilized in this section are introduced along with the presentation of particular equations. Other assumptions, more general in nature, will be noted before proceeding with the actual derivation of the equations.

It is assumed that the ideal gas law holds for the reactants and products at the initial and final conditions, respectively. This is equivalent to the assumption that the actual temperatures and pressures of the initial and final mixtures in all cases are in a region where the use of the ideal gas law results in negligible error. For the final mixture, this is a reasonable assumption in all cases in view of the extremely high temperatures occuring in detonation. For the initial mixture, this is a reasonable assumption if the initial conditions are sufficiently far from the critical conditions of pressure and temperature for each component so that compressibility is a negligible factor. Thus, in all cases where an equation of state has been required for describing the states of the initial and final mixtures at the initial and final conditions, respectively, the ideal gas law has been used.

The law of additives, which is valid for mixtures of ideal gases, is used to calculate the values of molecular weight  $(M_w)$  and molal heat capacity  $(C_p)$  for the initial and final mixtures.

It is assumed that chemical equilibrium is attained in the detonation wave.

The following equations of conservation across the detonation front are applicable:

Conservation of mass:

$$\rho_1 \quad V_1 = \rho_2 \quad V_2 \tag{1}$$

Conservation of momentum (considering only normal entry into the front):

$$P_1 + \rho_1 V_1^2 = P_2 + \rho_2 V_2^2$$
 (2)

Conservation of energy:

$$Je_1 + \frac{V_1^2}{2} + \frac{P_1}{\rho_1} + Jh = Je_2 + \frac{V_2^2}{2} + \frac{P_2}{\rho_2}$$
 (3)

The velocity of sound in the initial and final mixtures, respectively, may be expressed in the form

$$a_1^2 = \sqrt[p]{\frac{P_1}{\rho_1}} , \qquad (4)$$

$$a_2^2 = \sqrt[2]{\frac{P_2}{P_2}}$$
 (5)

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The assumption is made here that the adiabatic law holds for infinitesimal changes in the initial mixture at the initial conditions and in the final mixture at the final conditions. This is merely the usual assumption made in deriving the expression for the speed of sound.

Since  $M_1 = V_1/a_1$  and  $M_2 = V_2/a_2$ , then with equations (4) and (5),

$$V_1^2 = M_1^2 a_1^2 = \frac{Y_1 M_1^2 P_1}{\rho}$$
, (6)

$$V_2^2 * M_2^2 a_2^2 * \frac{\nabla_e M_2^2 P_2}{\rho_e}$$
 (7)

By definition, the energy release function is

$$B = \frac{Jh}{a_1}$$
 (8)

From equations (8) and (4),

$$Jh = Ba_1^2 = \frac{B \chi_1 P_1}{\rho_1} . \tag{9}$$

From the conservation of momentum, equation (2) and equations (6) and (7), the pressure ratio across the detonation front may be expressed as

$$\frac{P_2}{P_1} = \frac{\gamma_1 M_1^2 + 1}{\gamma_2 M_2^2 + 1}$$
 (10)

From the conservation of mass equation (1) and equations (6) and (7),

$$\frac{\rho_2}{\rho_1} = \left(\frac{M_1^2}{M_2^2}\right) \left(\frac{\zeta_1}{\zeta_2}\right) \left(\frac{\rho_1}{\rho_2}\right) . \tag{11}$$

and from equations (10) and (11), the density ratio may be expressed in the form

$$\frac{\rho_{z}}{\rho_{i}} = \left(\frac{\gamma_{i} M_{i}^{2}}{\gamma_{i} M_{i}^{2}+i}\right) \left(\frac{\gamma_{z} M_{z}^{2}+i}{\gamma_{z} M_{z}^{2}}\right) \qquad (12)$$

From equations (12) and (10),

$$\frac{P_2 P_1}{P_1 P_2} = \frac{V_2 M_2^2}{(V_2 M_2^2 + i)^2} \cdot \frac{(V_1 M_1^2 + i)^2}{V_1 M_1^2} . \tag{13}$$

The equation of state for ideal gases may be expressed in the following form:

$$\frac{P}{\rho} = \frac{JRT}{M_W} , \qquad (14)$$

Applying equation (14) to the state conditions before and after the detonation wave, designated by subscripts 1 and 2, respectively, the following relationship is obtained:

$$\frac{P_2 \rho_i}{P_1 \rho_2} = \frac{T_2 M_{w_1}}{T_1 M_{w_2}}.$$
 (15)

The dimensionless quantity  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$  is referred to hereafter as "the temperature-mole-

cular weight ratio. " From equations (13) and (15), the following expression for the temperature-molecular weight ratio is obtained:

$$\frac{T_{2} M_{W_{1}}}{T_{1} M_{W_{2}}} = \frac{\mathcal{T}_{2} M_{2}^{2}}{(\mathcal{T}_{2} M_{2}^{2} + 1)^{2}} \cdot \frac{(\mathcal{T}_{1} M_{1}^{2} + 1)^{2}}{\mathcal{T}_{1} M_{1}^{2}} . \tag{16}$$

The mass and momentum conservation equations ((1) and (2)) have been utilized to obtain non-dimensional expressions for  $\frac{P_2}{P_1}$ ,  $\frac{\rho_2}{\rho_i}$  and  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$  (equations (10), (12) and (16)).

In order to similarly transform the conservation of energy equation (3), consideration of the internal energy term is required.

The first law of thermodynamics can be expressed as

$$dq = de + \frac{1}{J} dw. ag{17}$$

For the case where the only force acting on the system is a fluid pressure, dw = Pdv, and equation (17) becomes

$$dq = de + \frac{1}{T} P dv . (18)$$

Differentiating equation (18) with respect to T, holding v constant, gives the following relationship:

$$\left(\frac{\delta q}{\delta T}\right)_{V} = \left(\frac{\delta e}{\delta T}\right)_{V} . \tag{19}$$

By definition, the specific heat at constant volume is

$$C_{V} = \left(\frac{\delta q}{\delta T}\right)_{V}, \qquad (20)$$

and it follows from equation (19) that

$$c_{v} = \left(\frac{\partial f}{\partial e}\right)_{v} . \tag{21}$$

For ideal gases,  $c_v$  is independent of pressure (or volume) and therefore is a function of temperature only. Hence, the following expression may be obtained from equation (21):

$$c_{v} = \frac{de}{dT}, \tag{22}$$

or 
$$de = c_v dT, (23)$$

or 
$$\Delta e = \int c_V dT$$
. (24)

For the internal energy per unit mass before and after passage through the detonation wave, respectively, equation (24) becomes

$$e_1 - (e_0)_{initial \ mixture} = \int_{T_0}^{T_1} (c_v)_{initial \ mixture} dT$$
, (25)

e<sub>2</sub> - (e<sub>0</sub>)<sub>final mixture</sub> = 
$$\int_{T_0}^{T_2} (c_v)_{final mixture} dT$$
, (26)

where subscripts 1 and 2 refer to the conditions before and after the detonation wave, respectively, and subscript "o" refers to the reference temperature  $T_o$ . Now, introducing  $\overline{c}_v$  as a mathematical convenience, equations (25) and (26) become

$$e_1 - (e_0)_{initial \ mixture} = \overline{c}_{v_{1,0}} (T_1 - T_0),$$
 (27)

$$e_2 - (e_0)_{\text{final mixture}} = \overline{c}_{V_2, o} (T_2 - T_0),$$
 (28)

where

e the mean specific heat at constant volume for the initial mixture between state 1 and the reference state (T<sub>0</sub>),

 $\overline{c}_{v_2, o}$  = the mean specific heat at constant volume for the final mixture between state 2 and the reference state  $(T_o)$ .

The reference temperature  $T_0$  will be taken to be absolute zero. In order to facilitate simplification of the generalized detonation equations, two assumptions will be made which the authors do not feel will appreciably affect the accuracy of the final generalized equations. It is assumed that the internal energy of the initial mixture at the reference temperature  $(T_0 = 0^\circ K)$  is equal to the internal energy of the final mixture at the reference temperature  $(T_0 = 0^\circ K)$ , i.e.,

The mean specific heats,  $\overline{c}_{v_1,o}$  and  $\overline{c}_{v_2,o}$ , will be replaced by the point properties  $c_{v_1}$  and  $c_{v_2}$ , respectively. This latter substitution or assumption will permit expression of the generalized detonation equations in terms only of the point properties of the initial and final mixtures, respectively.

Based on the foregoing statements, equations (27) and (28) may be combined to yield the following expression for the difference in internal energies across the detonation wave:

$$e_2 - e_1 * c_{v_2} T_2 - c_{v_1} T_1$$
 (29)

Utilizing equation (29), the conservation of energy equation (3) becomes

$$Jc_{v_1}T_1 + \frac{P_1}{\rho_1} + \frac{{v_1}^2}{2} + Jh = Jc_{v_2}T_2 + \frac{P_2}{\rho_2} + \frac{{v_2}^2}{2}$$
 (30)

For the case of ideal gases, from equation (14),

$$T_1 = \frac{M_{w_1}}{JR} \cdot \frac{P_1}{\rho_1} \text{ and } T_2 = \frac{M_{w_2}}{JR} \cdot \frac{P_2}{\rho_2}$$
 (31)

Using equations (31), the energy equation (30) may be written

$$\left(\frac{c_{v_1}M_{w_1}}{R} + 1\right)\frac{P_1}{\rho_1} + \frac{v_1^2}{2} + Jh = \left(\frac{c_{v_2}M_{w_2}}{R} + 1\right)\frac{P_2}{\rho_2} + \frac{v_2^2}{2}.$$
 (32)

Since  $c_{v_1}M_{w_1} = C_{v_1}$ , and  $c_{v_2}M_{w_2} = C_{v_2}$ , and since, for ideal gases,

 $R = C_{p_1} - C_{v_1} = C_{p_2} - C_{v_2}$ , equation (32) becomes

$$\left(C_{p_1}^{C_{p_1}}\right)^{\frac{p_1}{\rho_1}} + \frac{v_1^2}{2} + Jh = \left(C_{p_2}^{C_{p_2}}\right)^{\frac{p_2}{\rho_2}} + \frac{v_2^2}{2}, \qquad (33)$$

and, since  $\sum_{i=1}^{\infty} C_{p_1}/C_{v_1}$ , and  $\sum_{i=1}^{\infty} C_{p_2}/C_{v_2}$ , the following equation is obtained from equation (33):

$$\frac{Y_1}{Y_1-1} \frac{P_1}{P_1} + \frac{V_1^2}{2} + Jh = \frac{Y_2}{Y_2-1} \frac{P_2}{P_2} + \frac{V_2^2}{2}$$
 (34)

Equation (34) is the approximate energy equation utilized by D. G. Samaras in his paper on endothermic and exothermic discontinuities  $(1948)^{\frac{1}{2}}$  and involves approximations believed to be negligible, as discussed above.

From equations (6), (7), (9) and (34), the following equation is obtained:

$$\frac{\gamma_{1}}{\gamma_{-1}} + \frac{\gamma_{1} M_{1}^{2}}{2} + \gamma_{1} B = \left(\frac{\gamma_{2}}{\gamma_{2}-1} + \frac{\gamma_{2} M_{2}^{2}}{2}\right) \frac{P_{2} \rho_{1}}{P_{1} \rho_{2}}, \qquad (35)$$

and, from equations (35) and (13),

$$B = \left(\frac{\chi^{2}}{\chi^{2}-1} + \frac{\chi^{2} M_{5}^{2}}{2}\right) \cdot \frac{\chi^{2} M_{5}^{2}}{(\chi^{2} M_{5}^{2}+1)^{2}} \cdot \frac{(\chi^{1} M_{1}^{2}+1)^{2}}{\chi^{2} M_{1}^{2}} - \frac{1}{\chi^{1}-1} - \frac{M_{1}^{2}}{2} \cdot (36)$$

This is a form of the Hugoniot equation, in non-dimensional form, based on the energy equation (34).

The Chapman-Jouguet hypothesis that the product gases (final mixture) leave the detonation front at a velocity equal to the local speed of sound in those gases may be expressed in the form

$$M_2 = 1.0$$
 . (37)

With this substitution, equation (36) reduces to

$$B = \left(\frac{\gamma_{g}}{\gamma_{g-1}} + \frac{\gamma_{g}}{2}\right) \cdot \frac{\gamma_{g}}{(\gamma_{g+1})^{2}} \cdot \frac{(\gamma_{i} M_{i}^{2} + 1)^{2}}{\gamma_{i}^{2} M_{i}^{2}} - \frac{1}{\gamma_{i-1}} - \frac{M_{i}^{2}}{2}, \qquad (38)$$

which, upon rearrangement, gives the following simplified expression:

$$B = \frac{1}{V_a^2 - 1} \left[ \frac{M_i^2}{2} + \left( \frac{V_a}{V_i} \right)^2 \left( \frac{1}{2 M_i^2} + V_i \right) \right] - \frac{1}{V_i - 1}$$
 (39)

This is a non-dimensional form of the equation commonly referred to as the Chapman-Jouguet equation, based on the energy equation (34).

Applying the Chapman-Jouguet hypothesis (equation (37)) to equations (0), (12) and (16)), expressions for the pressure ratio, density ratio and temperature-molecular weight ratio across the detonation wave, in non-dimensional form, may be obtained as follows:

$$\frac{P_2}{P_1} = \frac{1}{V_2 + 1} \left( V_1 M_1^2 + 1 \right) , \qquad (40)$$

$$\frac{\rho_{2}}{\rho_{1}} = \frac{\sqrt[3]{2} + 1}{\sqrt[3]{2}} \left( \frac{\sqrt[3]{1} M_{1}^{2}}{\sqrt[3]{1} M_{1}^{2} + 1} \right) , \qquad (41)$$

$$\frac{T_2 M_{W_1}}{T_1 M_{W_2}} = \frac{\chi_2}{(\chi_2 + 1)^2} \cdot \frac{(\chi_1 M_1^2 + 1)^2}{\chi_1 M_1^2}.$$
 (42)

The above three equations, together with equation (39), are the non-dimensional generalized detonation equations which form the basis for calculations and graphical presentations in this report.

Equations (40), (41) and (42) are presented graphically in Figures 6, 7 and 8, respectively. In each of these three graphs,  $\mathcal{L}_i M_1^2$  has been made the abscissa. Figure 9 is a plot of  $M_1^2$  vs  $\mathcal{L}_i M_1^2$  with lines of constant  $\mathcal{L}_i$  and is presented merely to show the overall range of the term  $\mathcal{L}_i M_1^2$  encountered in the detonation of gaseous mixtures.

The interrelationship between the variables  $P_2/P_1$ ,  $\rho_2/\rho_1$  and  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$  (the dependent variables shown in equations (40), (41) and (42), respectively) is presented in

graphical form in Figure 10. It will be noted that this graph is simply a plot of equation (15), which is the equation of state for ideal gases.

### IIL REARRANGEMENT OF THE CHAPMAN-JOUGUET EQUATION

The Chapman-Jouguet equation, (39), may be rearranged and expressed in the following manner in order to facilitate the graphical presentation and to simplify the computational procedure:

$$M_1^2 + \left(\frac{\chi_2}{\chi_1}\right)^2 \frac{1}{M_1^2} = 2(\chi_2^2 - 1)B + 2\frac{\chi_2^2 - \chi_1}{\chi_1(\chi_1 - 1)}$$
 (43)

Equation (43) may be simplified by replacing the terms on the left hand side of the equation by a single symbol,  $\overline{M}_1^2$ , and replacing the two terms on the right hand side of the equation by simple functions of the dimensionless parameters E and  $\Gamma$ , respectively. The term  $\overline{M}$ , is introduced at this point and is given the name "pseudo-Mach number." The parameter E is hereafter referred to as "the dimensionless parameter epsilon," and the dimensionless parameter  $\Gamma$  is hereafter referred to as "the gamma function." Thus, equation (43) may be expressed in either of the following two forms:

$$\overline{M}_1^2 = 2 (\chi_2^2 - 1) B + 2 \frac{\chi_2^2 - \chi_1}{\chi_1 (\chi_1 - 1)}$$
, (44)

or

$$\overline{\mathbf{M}_{i}^{2}} = 2 \mathbf{R} + 2 \mathbf{\Gamma}_{j} \tag{45}$$

where

$$\overline{M}_{1}^{2} = M_{1}^{2} + \left(\frac{\overline{\chi}_{2}}{\chi_{1}}\right)^{2} \frac{1}{M_{1}^{2}}$$
, (46)

$$E = (\chi_2^2 - I) \quad B \quad , \tag{47}$$

$$\Gamma = \frac{\chi_2 - \chi_1}{\chi_1 (\chi_1 - 1)} . \tag{48}$$

It should be noted that equations (43), (44) and (45) are merely rearrangements of the non-dimensional form of the Chapman-Jouguet equation, based on the energy equation (34).

Equations (47), (48), (45) and (46) are presented graphically in Figures 1, 2, 3 and 4 respectively. These figures present the non-dimensional Chapman-Jouget equation in graphical form. Figure 5 is a plot of  $\sqrt[6]{t_1}$  vs  $\sqrt[6]{t_2}$  with lines of constant  $\sqrt[6]{t_1}$  and is presented merely to show the overall range of  $\sqrt[6]{t_2}/\sqrt[6]{t_1}$  encountered in the detonation of gaseous mixtures.

### IV. BASIC DATA REQUIRED

- A. Data on heats of formation for the reactants at the initial temperature and pressure, or these data at a standard temperature and pressure in conjunction with data on the change C<sub>p</sub> with temperature between the standard conditions and the initial conditions.
- B. Data on molal heat capacity,  $C_{p(i)}$ , for each component of the initial mixture at the initial conditions.
- C. Data on molal heat capacity,  $C_{p(j)}$ , for each component of the final mixture at the final conditions in the detonation wave.
- D. Data on equilibrium constants for dissociation reactions involving the components of the final mixture at the final conditions in the detonation wave.

### V. DISCUSSION OF THE CALCULATION PROCEDURE

The calculation procedure based on the generalized detonation equations derived in this report is presented in detail in Appendix I. Immediately preceding this procedure a complete outline of the steps involved is given in order to facilitate its use.

Indications are that previous methods for calculating detonation parameters have utilized directly the classical Chapman-Jouguet equations involving dimensional quantities, 6.12.13.14.15 whereas in this report the calculation procedure utilizes generalized equations and non-dimensional parameters throughout.

This method of calculation involves a process of successive approximations or a trial-and-error procedure. This is necessitated by the fact that a direct solution of the generalized equations is still impractical, if not impossible. The fact that dissociation of the product gases is taken into account in the calculations results in a triple trial-and-error procedure. This procedure involves the assignment of arbitrary values to final pressure, temperature and composition and then, by a method of successive approximations, the values of these quantities are adjusted so that the calculated values approach the assumed values. Dimensionless quantities and graphical presentation of the generalized data, as described above, are used to simplify and streamline considerably the calculation procedure. The method presented provides for a rapid convergence of the trial-and-error calculations by simple graphical techniques. The calculation procedure also provides for continual self-checking of results at intermediate points throughout the entire calculation, i.e., errors can be detected and discarded with little loss of time.

Considerable effort has been expended in the preparation of this method to insure that the information contained therein is clearly understandable to any research scientist who might have occasion to be working in this field (aeronautical, mechanical, chemical, electrical or propulsion engineer, or chemist). The calculation procedure, charts and graphs are presented so that even those not experienced in the detonation field can readily utilize the information presented to obtain accurate values of the required detonation parameters by routine manual calculation with a minimum of time, ingenuity, instruction and/or supervision. Where extensive data at different initial conditions and for various gaseous mixtures are desired, this method is also very readily adaptable to calculations by automatic computing machines.

### VI. DISCUSSION OF THE SAMPLE CALCULATION

A sample calculation utilizing the calculation procedure presented in this report is given in Appendix III. Calculations are presented for the determination of the detonation parameters for a stoichiometric mixture of hydrogen and oxygen at an initial temperature and pressure of 18°C and atmospheric pressure, respectively.

Results obtained in these calculations (taking into consideration atomic oxygen in the reaction products and utilizing consistent data throughout for specific heats and equilibrium constants in the range of the final temperature) indicated a detonation temperature and pressure of 3536°K and 17.2 atmospheres, respectively, and a detonation velocity of 8962 ft/sec. Results obtained by this calculation procedure neglecting the dissociation of molecular oxygen but still utilizing the same consistent data for specific heat and equilibrium constants gave values of the detonation temperature and pressure of 3593°K and 17.3 atmospheres, respectively, and a detonation velocity of 8991 ft/sec.\*

<sup>\*</sup>This calculation was performed by the authors, but detailed calculations are not presented in this report.

Calculated data obtained from literature sources indicate a detonation temperature and pressure of 3583°K and 18.05 atmospheres, respectively, and a detonation velocity of 9203 ft/sec, these values having been calculated neglecting the dissociation of molecular oxygen<sup>12,13</sup>. The difference between the calculated data presented in this report and that reported in the literature is probably due to (a) the use of slightly different specific heat and equilibrium constant data in calculating the detonation parameters, and (b) the simplifying assumptions utilized in deriving the generalized equations, which apparently did not affect greatly the accuracy of the final results.

Preliminary experimental data obtained by the authors in a one-inch diameter detonation tube showed very close agreement (approximately 1-1/2%) with the data calculated by the procedure reported. The experimental value obtained for the detonation velocity was 9100 ft/sec.

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### APPENDIX I

### CALCULATION PROCEDURE

Detailed Outline of Calculation Steps
Detailed Calculation Procedure

For convenience in using this calculation procedure, a detailed outline of the steps involved is presented below.

- Step 1. Proposed Scheme of Reaction
- Step 2. Setting-Up of the Partial Pressure Equations for Use in the Calculation of the Composition of the Final Mixture
- Step 3. Solution of the Simultaneous Partial Pressure (p<sub>j</sub>) Equations in General Terms
- Step 4. Fixing of the Initial Composition, ni'
- Step 5. Particular Solution of the Partial Pressure (p<sub>j</sub>) Equations for a Specific Initial Mixture Composition (n<sub>i</sub>')
- Step 6. Fixing of the Initial Pressure, p1
- Step 7. Fixing of the Initial Temperature, t
- Step 8. Calculation of the Properties of the Initial Mixture
- Step 9. Assumption of the Final Pressure, p,
- Step 10. Assumption of the Final Temperature, T2
- Step 11. Determination of Numerical Values of the Equilibrium Constants for the Dissociation Reactions
- Step 12. Calculation of the Partial Pressures (pi's) in the Final Mixture
- Step 13. Determination of the Total Number of Moles (n2) of the Final Mixture.
- Step 14. Calculation of the Composition (n;'s) of the Final Mixture
- Step 15. Calculation of the Properties of the Final Mixture
- Step 16. Calculation of the Heat of Reaction,  $\Delta H_R$
- Step 17. Calculation of the Energy Release Function, B
- Step 18. Determination of the Value of the Gamma Ratio,  $\chi_2/\chi_1$
- Step 19. Calculation of the Molecular Weight Ratio,  $M_{w_2}/M_{w_1}$
- Step 20. Determination of the Value of the Dimensionless Parameter  $\, {f E} \,$
- Step 21. Determination of the Value of the Gamma Function,  $\Gamma$

- Step 22. Determination of the Value of the Pseudo-Mach Number (Squared),  $\overline{M}_1^2$
- Step 23. Determination of the Value of the Detonative Mach Number (Squared), M.
- Step 24. Determination of the Value of  $\chi_{1}^{2}$
- Step 25. Determination of the Value of the Temperature-Molecular Weight Ratio,  $T_2M_{w_1}/T_1M_{w_2}$
- Step 26. Calculation of the Final Temperature, T2
- Step 27. Determination of the Value of the Pressure Ratio,  $p_2/p_1$
- Step 28. Calculation of the Final Pressure, p2
- Step 29. Calculation of the Detonative Mach Number, M,
- Step 30. Calculation of the Detonation Velocity, D
- Step 31. Calculation of Additional Data
  - a. Density ratio,  $\rho_2/\rho_1$ (or Volume Ratio,  $v_1/v_2$ )
  - b. Initial Density,  $\rho_i$  (or  $\frac{1}{v_1}$ )
  - c. Final Density,  $\rho_2$  (or  $\frac{1}{v_2}$ )

The detailed calculation procedure is presented below and follows the outline given directly above.

#### Step 1. Proposed Scheme of Reaction

Set down initially a proposed scheme of reaction, including all possible or probable products:

$$n_A^i A + n_B^i B + \dots \longrightarrow n_U^i U + n_V^i V + \dots,$$
 (49)

n' A and n' B being known quantities, n u and n being initially unknown quantities. The number of possible products included in this reaction scheme will depend upon experience and information available on the particular reaction being investigated.

The calculation procedure should be started on the basis of the above proposed reaction scheme. After the first calculation of composition, based on the first assumed temperature and pressure, it may be apparent that one or more of the products in the proposed reaction scheme is

present only in negligible amount. If this is the case, then it may be advisable to modify the above proposed scheme of reaction to eliminate the negligible constituents before proceeding with the second calculation of composition.

Step 2. Setting-Up of the Partial Pressure (p<sub>j</sub>) Equations for use in the Calculation of the Composition of the Final mixture

#### a. Dissociation Reactions and Equilibrium Equations

Set down dissociation reactions involving each of the products included in the above reaction scheme, along with the corresponding equilibrium equations expressed in terms of the partial pressures of the constituents. For example, if H<sub>2</sub>O, H<sub>2</sub> and O<sub>2</sub> are three of the components of the final mixture, then

$$H_2O \implies H_2 + 1/2 O_2$$
 (50)

and

$$\frac{p_0^{\frac{1}{2}} p_{H_2}}{p_{H_2^0}} = K_{p(H_2^0)} = \text{equilibrium constant.}$$
 (51)

In selecting the dissociation reactions to be used in these calculations, care must be excercised to insure that the equilibrium equations obtained are independent equations. Each successive dissociation reaction written must involve a new component of the final mixture if the corresponding equilibrium equation is to be independent of the other equilibrium equations. Another factor to be considered in the selection of the dissociation reactions to be used is the availability of equilibrium constant data for particular reactions.

#### b. Additional Equations Required

The dissociation reactions have provided a number of independent equations involving the partial pressure of each component of the final mixture. In order to solve for each of these partial pressures, the total number of independent equations available must equal the number of components in the final mixture. Independent equations required in addition to the equilibrium equations are obtained from (1) application of Dalton's Law, i.e., summation of partial pressures, and (2) application of the law of conservation of mass for each element involved in the proposed scheme of reaction.

#### c. Summation of Partial Pressures in the Final Mixture

Set down the equation for the summation of the partial pressures of the components of the final mixture:

$$p_U + p_V + p_W + \cdots = \sum_{j} p_j = p_2.$$
 (52)

#### d. Conservation of Mass for Each Element

Set down the mass-balance equation stating the conservation of atomic type for each chemical element involved in the proposed reaction scheme. The number of these mass-balance equations is equal to the number of chemical elements involved. The law of conservation of mass yields, for elements F and G, the following two equations:

For element F:

$$\pi_A^F N_A + \pi_B^F N_B + \dots = n_U^F N_U + n_V^F N_V + \dots,$$
 (53)

or

$$\sum_{i} n_{i}^{\dagger} N_{i} = \sum_{j} n_{j}^{\dagger} N_{j} . \qquad (54)$$

For element G:

$$n_A^e N_A + n_B^e N_B + \dots = n_U^e N_U + n_V^e N_V + \dots,$$
 (55)

or

$$\sum_{i} \pi_{i}^{G} N_{i} = \sum_{j} n_{j}^{G} N_{j}. \qquad (56)$$

Divide each mass balance equation by n<sub>2</sub> (value initially unknown) as shown below. For element F, the equation obtained is

$$\frac{1}{n_2} \sum_{i} n_i^{\dagger} N_i = \sum_{j} \frac{n_j}{n_2} N_j; \qquad (57)$$

and, similarly, for element G,

$$\frac{1}{n_2} \sum_{i} n_i^{i} N_i = \sum_{j} \frac{n_j}{n_2} N_j.$$
 (58)

Substitute into the mass balance equation for each element the following equality, obtained from Dalton's Law of Additive Pressures and the ideal-gas law:

$$\frac{\mathbf{n_j}}{\mathbf{n_2}} = \frac{\mathbf{P_j}}{\mathbf{P_2}} \tag{59}$$

This yields a mass balance equation, for each element, in terms of the partial pressures  $(p_j's)$  of components of the final mixture. For element F, this yields the equation

$$\frac{1}{n_2} \sum_{i} n_i^{F} N_i = \sum_{j} \frac{p_{j}^{F}}{p_2} N_j , \qquad (60)$$

and, for element G, the equation

$$\frac{1}{n_2} \sum_{i} n_i^{\alpha} N_i = \sum_{j} \frac{p_j}{p_2} N_j ; \qquad (61)$$

or, rearranging slightly, for element F,

$$\frac{1}{n_2} \sum_{i} n_i^{\dagger} N_i = \frac{1}{p_2} \sum_{j} p_j^{\dagger} N_j , \qquad (62)$$

and, for element G,

$$\frac{1}{n_2} \sum_{i} n_i^{e} N_i = \frac{1}{p_2} \sum_{j} p_j^{e} N_j . \qquad (63)$$

As stated above, the number of these mass balance equations (equations (62) and (63)) is equal to the number of chemical elements involved in the proposed reaction scheme. Any two of the mass balance equations may be combined to yield an independent equation in terms of the partial pressures  $(p_j's)$ , by elimination of  $n_2$  and  $p_2$ . For example, for elements F and G, dividing equation (62) by (63) yields

$$\frac{\sum_{i}^{j} n_{i}^{i} N_{i}}{\sum_{i}^{j} n_{i}^{i} N_{i}} = \frac{\sum_{j}^{j} p_{j}^{j} N_{j}}{\sum_{j}^{j} p_{j}^{j} N_{j}}.$$
 (64)

If there are three chemical elements involved in the reaction scheme, an additional independent equation in terms of the partial pressures (p<sub>j</sub>'s) may be obtained by using a new combination of mass balance equations. If more than three elements are involved, additional independent equations in terms of the partial pressures may be obtained in the same manner.

Step 3. Solution of the Simultaneous Partial Pressure (p<sub>j</sub>) Equations in General Terms

Step 2 has provided a number of independent partial pressure equations equal to the number of components in the final mixture. Solve these equations simultaneously to obtain a general solution. When the number of products of reaction is small, this solution of the simultaneous equations may be accomplished by direct algebraic methods. When the number of reaction products is large, it may be desirable to utilize methods of solution such as those reported by Edse,  $\frac{8}{}$  Huff and Calvert,  $\frac{5}{}$  Huff and Morrell,  $\frac{9}{}$  Hershfelder,  $\frac{10}{}$  etc.

Particular solutions for specific initial mixtures ( $n_i$ ) are obtained from the general solution by fixing the initial mixture composition (Steps 4 and 5).

Step 4. Fixing of the Initial Mixture Composition, ri

Fix the initial mixture composition,  $n_i$ , at the value to be used throughout the remainder of the calculations.

Step 5. Particular Solution of the Partial Pressure (p<sub>j</sub>) Equations for a Specific Initial Mixture Composition (n'<sub>i</sub>)

Substitute the numerical value for the initial mixture composition  $(n_i)$ , assigned in Step 4, into the general solution obtained in Step 3, to obtain a particular solution of the partial pressure  $(p_j)$  equations for a specific initial mixture composition.

Step 6. Fixing of the Initial Pressure, p

Fix the initial pressure, p<sub>1</sub>, at the value to be used throughout the remainder of the calculations.

Step 7. Fixing of the Initial Temperature, t1

Fix the initial temperature, t<sub>1</sub>, at the value to be used throughout the remainder of the calculations.

Step 8. Calculation of the Properties of the Initial Mixture

a. Mole Fraction, x;

$$x'_{\mathbf{A}} = \frac{n'_{\mathbf{A}}}{n'_{\mathbf{A}} + n'_{\mathbf{B}} + \dots} , \qquad (65)$$

or

$$x_{i}^{i} = \frac{n_{i}^{i}}{\sum_{i}^{n_{i}^{i}}} = \frac{n_{i}^{i}}{n_{1}}$$
 (66)

b. Molecular Weight,  $M_{W_{1}}$ 

$$M_{w_1} = x_i^{\dagger} M_{w(A)} + x_i^{\dagger} B M_{w(B)} + \cdots$$

$$= \sum_{i} x_i^{\dagger} M_{w(i)} . \qquad (67)$$

c. Molal Heat Capacity at Constant Pressure,  $C_{p_1}$ 

$$C_{p_{1}} = x'_{A} C'_{p(A)} + x'_{B} C'_{p(B)} + \dots$$

$$= \sum_{i} x'_{i} C'_{p(i)} .$$
(68)

d. Specific Heat Ratio (or Molal Heat Capacity Ratio),  $\chi_{_{1}}$ 

$$\chi_{1} = \frac{C_{p_{1}}}{C_{v_{1}}} = \frac{C_{p_{1}}}{C_{p_{1}} - R} . \tag{69}$$

e. Heat of Formation\*, Δ H'f(i)

Calculate the heat of formation of each component of the initial mixture at the initial conditions.

$$\Delta H'_{f(i)} = \Delta H_{f_{L}(i)} - H'_{i} + \sum H'_{e(i)}$$
 (70)

If the component "i" is an element in its standard state of aggregation at  $t_1$ , then, by definition, it is apparent that its heat of formation  $\Delta H'_{f(i)}$  is equal to zero, e.g.,  $\Delta H'_{f(H_2)}$  and  $\Delta H'_{f(O_2)}$  are equal to zero when  $t_1$  is in a region where hydrogen and

<sup>\*</sup> In this report, the heat of formation is taken to be positive when heat is evolved. If enthalpy data for this calculation cannot be found in the literature, it can be calculated readily from available  $C_{\rm p}$  data.

oxygen are normally diatomic gases.

Step 9. Assumption of the Final Pressure, p2

Assume a value for the final pressure, p2.

Step 10. Assumption of the Final Temperature, T2

Assume a value for the final temperature, T2.

Step 11. Determination of the Numerical Values of the Equilibrium Constants for the Dissociation Reactions

Obtain the numerical values of the equilibrium constant for each of the dissociation reactions considered in Step 2, at the assumed temperature (T<sub>2</sub>), either from available data in the literature or by calculation.

Because of the trial-and-error nature of the overall calculation procedure, it is advisable to make a plot of equilibrium constant versus temperature for each of the dissociation reactions in the temperature range in which T<sub>2</sub> is expected to be. It is necessary that this be done at the very beginning of the calculation procedure in order to insure that consistent data are used and thereby facilitate rapid convergence of the trial-and-error calculations.

Step 12. Calculation of the Partial Pressures (pi's) in the Final Mixture

Using the equation(s) obtained in Step 5 in conjunction with the equilibrium equations set forth in Step 2, par. a, determine the numerical value of the partial pressure (p<sub>i</sub>) of each constituent of the final mixture.

Step 13. Determination of the Total Number of Moles (n2) of the Final Mixture

Set down the overall mass balance equation for the proposed scheme of reaction:

$$n'_A M_{w(A)} + n'_B M_{w(B)} + \dots = n_U M_{w(U)} + n_V M_{w(V)} + \dots$$
, (71)

or

$$\sum_{i} n_{i}^{\prime} M_{w(i)} = \sum_{j} n_{j} M_{w(j)}. \qquad (72)$$

Substitute in equation (72) the relationship

$$n_j = \frac{p_j}{p_2} n_2$$
, (73)

obtained from equation (59); then

$$\sum_{i} n_{i}^{i} M_{w(i)} = \frac{n_{2}}{p_{2}} \sum_{j} p_{j} M_{w(j)} , \qquad (74)$$

and

$$n_2 = p_2 \frac{\sum_{i}^{n_i} M_{w(i)}}{\sum_{j}^{p_j} M_{w(j)}}$$
 (75)

This gives the total number of moles of final mixture corresponding to the number of moles of initial mixture shown in the proposed scheme of reaction (i.e.,  $n'_A + n'_B + \dots$ ).

#### Step 14. Calculation of the Composition (ni's)

Again using the relationship

$$n_{j} = \frac{p_{j}}{p_{2}} n_{2}$$
, (73)

obtained from equation (59), calculate the number of moles (n<sub>j</sub>) of each component in the final mixture.

#### Step 15. Calculation of the Properties of the Final Mixture

a. Mole Fraction, xi

$$x_{U} = \frac{n_{U}}{n_{2}} = \frac{p_{U}}{p_{2}}$$
, (76)

or

$$x_j = \frac{n_j}{n_2} = \frac{p_j}{p_2}$$
 (77)

b. Molecular Weight, Mw2

$$M_{w_2} = x_U M_{w(U)} + x_V M_{w(V)} + \dots = \sum_{j} x_{j} M_{w(j)}$$
 (78)

c. Molal Heat Capacity at Constant Pressure, Cp2

Determine the value of  $C_{p_2}$  at the above assumed value of  $T_2$  (Step 10), using values of  $C_{p(j)}$  obtained from the literature.

$$C_{p_2} = x_U C_{p(U)} + x_V C_{p(V)} + \dots = \sum_{j} x_j C_{p(j)}$$
 (79)

d. Specific Heat Ratio (or Molal Heat Capacity Ratio),  $\chi_2$ 

$$\chi_{2} = \frac{C_{p_{2}}}{C_{v_{2}}} = \frac{C_{p_{2}}}{C_{p_{2}} - R} . \tag{80}$$

If  $C_p$  data cannot be found in the literature, it can be readily calculated from available enthalpy data.

e. Heat of Formation\*,  $\Delta H_{f(j)}$ 

Calculate the heat of formation of each component of the final mixture at the initial conditions.

$$\Delta H'_{f(j)} = \Delta H_{f_{T}(j)} - H'_{j} + \sum H'_{e(j)}$$
 (81)

Step 16. Calculation of the Heat of Reaction, Energy Release per Mole of Initial Mixture ( $\Delta H_R$ )

$$\Delta H_{R} = \frac{n_{U} \Delta H'_{f(U)} + n_{V} \Delta H'_{f(V)} + \dots - n'_{A} \Delta H'_{f(A)} - n'_{B} \Delta H'_{f(B)} - \dots}{n_{A} + n_{B} + \dots},$$
(82)

or

$$\Delta H_{R} = \frac{\sum_{j}^{n_{j}} \Delta H'_{f(j)} - \sum_{i}^{n_{i}} \Delta H'_{f(i)}}{n_{1}}$$
 (83)

Step 17. Calculation of the Energy Release Function, B

By equations (8), (4) and (31) and by definition of "h" and  $\Delta H_R$ ,

<sup>\*</sup> See footnote in Step 8, par. e.

$$B = \frac{Jh}{a_1^2} = \frac{M_{w_1}h}{V_i RT_1} = \frac{\Delta H_R}{V_i RT_1}.$$
 (84)

Step 18. Determination of the Value of the Gamma Ratio,

Calculate this ratio from values of  $b_1$  and  $b_2$  or obtain from Figure 5.

Step 19. Calculation of the Molecular Weight Ratio,  $\frac{M_{w_2}}{M_{w_1}}$ 

Calculate this ratio from values of  $M_{W_1}$  and  $M_{W_2}$ .

Step 20. Determination of the Value of the Dimensionless Parameter E

Calculate E from values of Y2 and B or obtain from Figure 1.

$$\mathbf{E} = (\chi_2^2 - 1) \quad \mathbf{B} \tag{47}$$

Step 21. Determination of the Value of the Gamma Function,  $\Gamma$ 

Calculate  $\Gamma$  from values of  $\mathcal{X}_1$  and  $\mathcal{X}_2$  or obtain from Figure 2.

$$\Gamma = \frac{\chi_2^2 - \chi_i}{\chi_i (\chi_i - 1)} \tag{48}$$

Step 22. Determination of the Value of Pseudo-Mach Number (Squared),  $\overline{M}_1^2$ 

Calculate  $\overline{M}_1^2$  from values of E and  $\Gamma$  or obtain from Figure 3.

$$\overline{\mathbf{M}}_{1}^{2} = 2\mathbf{E} + 2\mathbf{\Gamma}, \tag{45}$$

or

$$\overline{M_1}^2 = 2 \left( \gamma_2^2 - 1 \right) B + 2 \frac{\gamma_2^2 - \gamma_1}{\gamma_1 (\gamma_1 - 1)} . \tag{44}$$

Step 23. Determination of the Value of the Detonative Mach Number (Squared),  $M_1^2$ .

Obtain the value of  $M_1^2$  from Figure 4.

$$\overline{M}_{1}^{2} = M_{1}^{2} + \frac{(\sqrt[3]{x_{1}})^{2}}{M_{1}^{2}}$$
 (46)

Step 24. Determination of the Value of  $V_1 M_1^2$ 

Calculate the value of  $V_i M_i^2$  from values of  $V_i$  and  $M_i^2$  or obtain from Figure 9.

Step 25. Determination of the value of the Temperature-Molecular Weight

Ratio, 
$$\frac{T_2M_{W_1}}{T_1M_{W_2}}$$

Calculate the value of  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$  from values of  $\chi_i^2$  and  $\chi_i^2$  or obtain from Figure 8. From equation (42),

$$\frac{T_2 M_{W_1}}{T_1 M_{W_2}} = \frac{\tilde{\chi}_2}{(\tilde{\chi}_2 + 1)^2} \left( \tilde{\chi}_1 M_1^2 + \frac{1}{\tilde{\chi}_1 M_1^2} + 2 \right). \tag{85}$$

Step 26. Calculation of the Final Temperature, T2

Calculate the value of  $T_2$  from values of  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$ ,  $\frac{M_{w_2}}{M_{w_1}}$  and  $T_1$ .

$$T_2 = \frac{T_2 M_{w_1}}{T_1 M_{w_2}} \cdot \frac{M_{w_2}}{M_{w_1}} \cdot T_1 \qquad (86)$$

If this calculated value of  $T_2$  does not check the assumed  $T_2$  (Step 10), re-assume  $T_2$  and repeat Steps 10 through 26. In order to facilitate rapid convergence between calculated and assumed values, the following procedure may be used. After two calculations of  $T_2$ , plot calculated  $T_2$  versus assumed  $T_2$  and draw a line through these two points, intersecting the 45° line. The value of  $T_2$  at the point of intersection usually is a very good indication of the correct value of  $T_2$  corresponding to the assumed  $T_2$ . This value of  $T_2$  should be used as the assumed  $T_2$  for third trial. A typical plot is shown in Figure 11.

Step 27. Determination of the Value of the Pressure Ratio,  $\frac{P_2}{P_1}$ 

Calculate the value of  $\frac{P_2}{P_1}$  using the values of  $\sqrt[8]{g}$  and  $\sqrt[8]{g}$  corresponding to the final checked value of  $T_2$  obtained in the trial-and-error calculations of Steps 10 through 26, or obtain from Figure 6. From

equation (40), since  $\frac{P_2}{P_1}$  is numerically equal to  $\frac{p_2}{P_1}$ ,

$$\frac{p_2}{p_1} = \frac{1}{|Y_2|+1} (|Y_1|M_1^2+1). \tag{87}$$

Step 28. Calculation of the Final Pressure, p2

$$p_2 = \frac{p_2}{p_1} \cdot p_1$$
 (88)

If this calculated value of  $p_2$  does not check the assumed  $p_2$  (Step 9), reassume  $p_2$  and repeat Steps 9 through 28 until the calculated and assumed values agree.

The graphical procedure outlined under Step 26 may be used to facilitate rapid convergence between calculated and assumed values of p<sub>2</sub>.

Step 29. Calculation of the Detonative Mach Number, M1

Calculate  $M_1$  from the value of  $M_1^2$  (Step 23) corresponding to the final checked values of  $p_2$  and  $T_2$  (Steps 26 through 28).

Step 30. Calculation of the Detonation Velocity, D

Since  $D = V_1 = M_{1}a_1$ , then with equations (4) and (14),

$$D = M_1 \sqrt{\frac{J \, \mathcal{K}_i \, R \, T_i}{M_{W_i}}} \quad . \tag{89}$$

Step 31. Calculation of Additional Data

a. Density Ratio,  $\frac{\rho_2}{\rho_1}$  (or Volume Ratio,  $\frac{v_1}{v_2}$ )

Calculate  $\frac{p_2}{p_1}$  (or  $\frac{v_1}{v_2}$ ) from the final checked values of  $\frac{p_2}{p_1}$ 

and  $\frac{T_2M_{w_1}}{T_1M_{w_2}}$  obtained in the trial-and-error calculations of Steps 26

through 28, or obtain from Figure 10.

$$\frac{\rho_{2}}{\rho_{1}} = \frac{P_{2}}{P_{1}} / \frac{T_{2}M_{W_{1}}}{T_{1}M_{W_{2}}} . \qquad (15)$$

b. Initial Density,  $\rho_1$  (or  $\frac{1}{v_1}$ )

By equation (14), for the initial mixture,

$$\rho_{\rm i} = \frac{M_{\rm w_1}P_1}{\rm JRT_1} \quad . \tag{90}$$

c. Final Density, 
$$\rho_2$$
 (or  $\frac{1}{v_2}$ )

$$\rho_2 = \frac{\rho_2}{\rho_1} \cdot \rho_1 \quad . \tag{91}$$

#### APPENDIX II

GRAPHS OF BASIC DATA-EQUILIBRIUM CONSTANTS AND
MOLAL HEAT CAPACITIES

As pointed out in the Calculation Procedure in this report, it is desirable that consistent data on equilibrium constants and molal heat capacities be available. Since the sample calculation presented in this report (Appendix III) is for a stoichiometric hydrogen-oxygen system, the necessary equilibrium constant data and molal heat capacity data for this system are presented in this appendix.

Values of the equilibrium constants ( $K_p$ 's) as a function of the absolute temperature for  $H_2O$ ,  $H_2$  and  $O_2$  are presented graphically in Figures 12 through 15, respectively.

Values of the molal heat capacities ( $C_p$ 's) as a function of absolute temperature for  $H_2O$ ,  $H_2$ ,  $O_2$  and OH are presented graphically in Figures 16 through 19, respectively.

The range of temperature in all these figures extends from 2000° to 5000°K, and it is expected that this range will cover all cases to be encountered in the detonation of gaseous mixtures.

In each of the figures described above, a smooth curve was drawn through the data plotted in order to insure that consistent data would be available for the trial-and-error calculations.

It is believed that these graphs of basic data described above will find practical use in other applications, particularly in other combustion calculations where trial-and-procedures are involved and where consistent thermodynamic data are essential.

# APPENDIX III SAMPLE CALCULATION

# CALCULATION OF DETONATIVE MACH NUMBER AND ASSOCIATED PARAMETERS FOR $\mathrm{H_2}$ - $\mathrm{O_2}$ MIXTURES

The step numbers used below are the same as those given in the calculation procedure outlined in Appendix I.

#### Step 1. Proposed Scheme of Reaction

$$n'_{A}A + n'_{B}B \longrightarrow n_{U}U + n_{V}V + n_{W}W + n_{X}X + n_{Y}Y + n_{Z}Z$$
 (92)

or, specifically. for H<sub>2</sub> - O<sub>2</sub> mixtures,

$$n_{H_2}^{I}H_2^{+} n_{0_2}^{I}O_2^{---}N_{H_{20}}H_2^{0} + n_{H_2}^{H_2}H_2^{-} + n_{0_2}^{0}O_2^{-} + n_{OH}^{OH} + n_{H}^{H} + n_{0}^{O}^{-}$$
 (93)

#### Step 2. Setting-Up of the Partial Pressure (pi) Equations

#### a. Dissociation Reactions

(1) 
$$H_2O \implies H_2 + \frac{1}{2}O_2$$
 (94)

$$K_{P(H_{20})} = \frac{P_{H_2} \cdot P_{02}^{\frac{1}{2}}}{P_{H_{20}}} . \tag{95}$$

(2) 
$$H_2O \implies \frac{1}{2}H_2 + OH$$
 (96)

$$K_{P(H_2O)}^{'} = \frac{P_{H_2}^{\frac{1}{2}} \cdot P_{OH}}{P_{H_2O}}$$
 (97)

$$(3) \quad \frac{1}{2} H_2 \Longrightarrow H \tag{98}$$

$$K_{p_{(H_{\overline{z}})}} = \frac{p_{H}}{p_{H_{\overline{z}}}^{\frac{1}{2}}} . \tag{99}$$

$$(4) \quad \frac{1}{2}O_2 \Longrightarrow O \tag{100}$$

$$K_{p_{(o_2)}} = \frac{p_0}{p_0^{\frac{1}{2}}} . \tag{101}$$

#### b. Summation of the Partial Pressures in the Final Mixture

$$p_U + p_V + p_W + p_X + p_Y + p_Z = \sum_j p_j * p_2$$
, (102)

or, specifically, for H<sub>2</sub> - O<sub>2</sub> mixtures,

$$P_{H_2O} + P_{H_2} + P_{O_2} + P_{OH} + P_{H} + P_{O} = P_{2}$$
 (103)

c. Conservation of Mass for Each Element

Using equation (64),

$$\frac{\sum_{i} n'_{i} N_{i}}{\sum_{j} n'_{i} N_{i}} = \frac{\sum_{j} p_{j} N_{j}}{\sum_{j} p_{j} N_{j}}, \qquad (64)$$

and taking elements F and G to be O2 and H2, respectively, then

$$\frac{n_{0_2}^i}{n_{H_2}^i} = \frac{p_{H_20}^i + 2p_{0_2}^i + p_{0H}^i + p_0^i}{2p_{H_20}^i + 2p_{H_2}^i + p_{0H}^i + p_H^i},$$
 (104)

or

$$\left(2 \frac{n_{0z}^{\prime}}{n_{Hz}^{\prime}} - I\right) p_{Hz^{0}} + 2 \frac{n_{0z}^{\prime}}{n_{Hz}^{\prime}} p_{Hz} + \frac{n_{0z}^{\prime}}{n_{Hz}^{\prime}} p_{H} + \left(\frac{n_{0z}^{\prime}}{n_{Hz}^{\prime}} - I\right) p_{OH} - 2p_{Oz} - p_{O} = 0. \quad (105)$$

## Step 3. Solution of the Simultaneous Partial Pressure (p<sub>j</sub>) Equations in General Terms

Step 2 has provided the following six independent partial pressure equations, which are sufficient to determine the partial pressures of each of the six components in the final mixture.

$$K_{p_{(H_20)}} = \frac{p_{H_2} p_{0_2}^{\frac{1}{2}}}{p_{H_20}},$$
 (95)

$$K_{p_{(H_2O)}}^{\prime} = \frac{p_{H_2}^{\frac{1}{2}} p_{OH}}{p_{H_2O}},$$
 (97)

$$K_{\rho_{(H_Z)}} = \frac{\rho_H}{\rho_{H_Z}} , \qquad (99)$$

$$Kp_{(o_2)} = \frac{p_o}{p_{o_2}}$$
, (101)

$$P_{H_2O} + P_{H_2} + P_{O_2} + P_{OH} + P_{H} + P_{O} = P_2,$$
 (103)

$$\left(2\frac{n_{Q_2}^i}{n_{H_2}^i}-i\right)p_{H_2Q}+2\frac{n_{Q_2}^i}{n_{H_2}^i}p_{H_2}+\frac{n_{Q_2}^i}{n_{H_2}^i}p_{H}+\left(\frac{n_{Q_2}^i}{n_{H_2}^i}-i\right)p_{QH}-2p_{Q_2}-p_{Q_2}-p_{Q_2}=0. \quad (105)$$

Rearranging equations (95), (97), (99) and (101),

$$P_{02} = \left(\frac{K_{p_{(H_20)}} P_{H_20}}{P_{H_2}}\right)^2,$$
 (106)

$$P_{OH} = \frac{K_{p_{(H_2O)}} p_{H_2O}}{p_{H_2}^{\frac{1}{2}}},$$
 (107)

$$P_{H} = K_{P_{(H_{2})}} P_{H_{2}}^{\frac{1}{2}},$$
 (108)

$$P_0 = K_{p_{(0_2)}} P_{0_2}^{\frac{1}{2}} = \frac{K_{p_{(0_2)}} K_{p_{(H_20)}} P_{H_20}}{P_{H_2}},$$
 (109)

and substituting these relationships into equations (103) and (105), respectively, the following two independent quadratic equations in terms of two variables,  $p_{H_2O}$  and  $p_{H_2}$ , are obtained:

$$\left(\frac{Kp_{(Hg0)}}{p_{Hg}}\right)^{2}p_{Hg0}^{2} + \left(\frac{K_{p_{(Hg0)}}^{1}}{p_{Hg}^{2}} + \frac{Kp_{(0g)}Kp_{(Hg0)}}{p_{Hg}} + i\right)p_{Hg0} + \left(p_{Hg} + Kp_{(Hg)}p_{Hg}^{2} - p_{g}\right) = 0, \quad (110)$$

$$2\left(\frac{K_{p_{(H_{2}0)}}}{p_{H_{2}}}\right)^{2}p_{H_{2}0}^{2} + \left[\frac{K_{p_{(D_{2})}}K_{p_{(H_{2}0)}}}{p_{H_{2}}} - \left(\frac{n_{O_{2}}^{\prime}}{n_{H_{2}}^{\prime}} - i\right)\frac{K_{p_{(H_{2}0)}}^{\prime}}{p_{H_{2}^{\frac{\prime}{2}}}} - 2\frac{n_{O_{2}}^{\prime}}{n_{H_{2}}^{\prime}} + i\right]p_{H_{2}0}$$

$$-\frac{n_{O_{2}}^{\prime}}{n_{H_{2}}^{\prime}}\left(2p_{H_{2}}^{\prime} + K_{p_{(H_{2})}}p_{H_{2}}^{\prime}\right) = 0.$$
(111)

#### Step 4. Initial Composition

For the purpose of this sample calculation, a stoichiometric mixture of hydrogen and oxygen will be considered, the initial composition being given by

$$n_{H_2}^{i} H_2 + n_{O_2}^{i} O_2$$
, (112)  
 $n_{H_2}^{i} = i$  and  $n_{O_2}^{i} = \frac{i}{2}$ ,

where

or  $H_2 + \frac{1}{2}O_2$ .

#### Step 5. Particular Solution of the Partial Pressure (pi) Equations

From the values for initial composition given in Step 4,

$$\frac{n'_{0_2}}{n'_{H_2}} = \frac{1}{2} . (113)$$

Substitution of the numerical value of this ratio into the general partial pressure equations (Equations (110) and (111)) obtained in Step 3 yields the following particular solution for a stoichiometric mixture of hydrogen and oxygen:

$$\left(\frac{K_{p_{(H_{2}0)}}}{p_{H_{2}}}\right)^{2} p_{H_{2}0}^{2} + \left(\frac{K_{p_{(0_{2})}} K_{p_{(H_{2}0)}}}{p_{H_{2}}} + \frac{K_{p_{(H_{2}0)}}^{1}}{p_{H_{2}}^{\frac{1}{2}}} + 1\right) p_{H_{2}0} + \left(p_{H_{2}} + K_{p_{(H_{2})}} p_{H_{2}}^{\frac{1}{2}} - p_{2}\right) = 0$$
(114)

and

$$2\left(\frac{K_{p_{(H_{2}0)}}}{p_{H_{2}}}\right)^{2} p_{H_{2}0}^{2} + \left[\frac{K_{p_{(0_{2})}} K_{p_{(H_{2}0)}}}{p_{H_{2}}} + \frac{K_{p_{(H_{2}0)}}^{1}}{2 p_{H_{2}}^{\frac{1}{2}}}\right] p_{H_{2}0} - \left(p_{H_{2}} + \frac{1}{2}K_{p_{(H_{2})}} p_{H_{2}}^{\frac{1}{2}}\right) = 0.$$
(115)

#### Step 6. Initial Pressure

For the purpose of this sample calculation, the initial pressure is fixed at 1 atmos., i.e.,

$$p_1 = 1$$
 atmos.

and

$$P_1 = 1.0133 \times 10^6 \text{ dynes/cm}^2$$
.

#### Step 7. Initial temperature

For the purpose of this sample calculation, the initial temperature is fixed at 18°C, i.e.,

$$t_1 = 18$$
°C.

and

$$T_1 = 291^{\circ}K.$$

#### Step 8. Properties of the Initial Mixture

#### a. Mole Fractions

$$X'_{H_{2}} = \frac{n'_{H_{2}}}{n'_{H_{2}} + n'_{O_{2}}} = \frac{1}{1 + \frac{1}{2}} = \frac{2}{3} ;$$

$$X'_{O_{2}} = \frac{n'_{O_{2}}}{n'_{H_{2}} + n'_{O_{2}}} = \frac{\frac{1}{2}}{1 + \frac{1}{2}} = \frac{1}{3} ;$$

$$\sum_{i} X_{i} = X'_{H_{2}} + X'_{O_{2}} = 1 .$$

#### b. Molecular Weight

$$M_{W_1} = \chi_{H_2}^{b} M_{W_{(H_2)}} + \chi_{O_2}^{l} M_{W_{(O_2)}}$$
  
=  $(\frac{2}{3})$  (2.0160) +  $(\frac{1}{3})$  (32.000)  
= 12.0107.

#### c. Molal Heat Capacity

At 
$$t_1 = 18^{\circ}C.$$
,  
 $C'_{P_{(Hg)}} = 6.88^{\frac{1}{2}} \text{ cal/(gm mole)(°K);}$ 

$$C'_{P_{(o_2)}} = 7.01 \frac{1}{cal/(gm mole)} (^{\circ}K);$$
 $C_{P_i} = X'_{H_2} C'_{P_{(H_2)}} + X'_{o_2} C'_{P_{(o_2)}}$ 
 $= (2/3)(6.88) + (1/3)(7.01) = 6.9233.$ 

d. Specific Heat Ratio

$$\overline{\delta}_{1} = \frac{C_{p_{1}}}{C_{p_{1}} - R} = \frac{6.9233}{6.9233 - 1.98719} = 1.4026.$$

e. Heat of Formation  $\Delta H'_{f(i)}$  of Each Reactant at the Initial Conditions

By definition of  $\Delta H'_{f(i)}$ , the heat of formation of an element (in its standard state of aggregation) at  $t_1$  is equal to zero.

Therefore,

$$\Delta H_{f_{(0_2)}}^i = 0.0;$$
 $\Delta H_{f_{(0_2)}}^i = 0.0.$ 

Step 9. Initial Assumption for p

$$p_2 = 20.00$$
 atmos.

Step 10. Initial Assumption for T2

$$T_2 = 3800^{\circ} K.$$

Step 11. Determination of Numerical Values of the Equilibrium Constants for the Dissociation Reactions

From Figures 12 through 15,

$$K_{p_{(H_20)}} = .3758 \text{ (atmos)}^{1/2} \text{ at } T_2 = 3800^{\circ} \text{K.}, *$$
 $K_{p_{(H_20)}}^{i} = .5984 \text{ (atmos)}^{1/2} \text{ at } T_2 = 3800^{\circ} \text{K.}, **$ 

<sup>\*</sup> This value also obtainable from Ref. 8, Table 4, by graphical interpolation.

<sup>\*\*</sup> This value also given in Ref. 7, page 25.

Step 12. Calculation of the Partial Pressures (pj's) in the Final Mixture

Referring to equations (114) and (115), let

$$c_{II} = \left(\frac{K_{p_{(HgO)}}}{p_{Hg}}\right)^2, \qquad (116)$$

$$c_{12} = \left(\frac{Kp_{(0_2)} Kp_{(H_20)}}{p_{H_2}} + \frac{K_{p_{(H_20)}}^{\prime}}{p_{H_2}} + 1\right), \qquad (117)$$

$$c_{13} = (p_{H_2} + K_{p(H_2)} p_{H_2}^{\frac{1}{2}} - p_2),$$
 (118)

$$k_{II} = 2 \left( \frac{K_{p_{(H_20)}}}{p_{H_0}} \right)^2 = 2 c_{II},$$
 (119)

$$k_{12} = \left(\frac{K\rho_{(0g)} K\rho_{(Hg0)}}{p_{Hg}} + \frac{K_{\rho_{(Hg0)}}^{1}}{2 p_{Hg}^{\frac{1}{2}}}\right),$$
 (120)

and

$$k_{13} = -(p_{H_2} + \frac{1}{2} K_{p_{(H_2)}} p_{H_2}^{\frac{1}{2}}).$$
 (121)

Then, from equation (114),

$$P_{Hg0} = \frac{-c_{1g} \pm \sqrt{c_{1g}^2 - 4c_{1l} c_{13}}}{2c_{1l}}, \qquad (122)$$

<sup>\*</sup> This value also given in Ref. 7, page 25.

<sup>\*\*</sup> This value also given in Ref. 6.

and, from equation (115),

$$P_{H_{20}} = \frac{-k_{12} \pm \sqrt{k_{12}^2 - 4k_{11}^2 k_{13}^2}}{2k_{11}}$$
 (123)

For the solution of equations (122) and (123), values of  $p_{He}$  have been assumed, as indicated in Table 1. These values of  $p_{H_2}$ , in conjunction with the values of the equilibrium constants (Step 11) and the assumed final pressure (Step 9) have been used to determine the values of the constants  $c_{11}$ ,  $c_{12}$ ,  $c_{13}$ , and  $k_{11}$ ,  $k_{12}$ ,  $k_{13}$  (Equations (116) through (121)). Equations (122) and (123) were then solved to obtain two values of  $P_{H=0}$ , indicated in Table 1 as P<sub>N20</sub> (Equation (122)) and P<sub>N20</sub> (Equation (123)). The difference between these two values of  $p_{H20}$  is designated in Table 1 as  $\Delta p_{H_{20}}$ . The values of  $\Delta p_{H_{20}}$  obtained were plotted as a function of the assumed  $\rho_{H_{2}}$  (Figure 20). After two trials (i.e., two assumed values of  $p_{H_2}$ ), the value of  $p_{H_2}$  at the point where the curve crossed the  $\Delta p_{H_20}$ line, was taken as the assumed  $p_{H_2}$  for the following trial. The fourth trial gave sufficiently accurate values of  $\boldsymbol{p}_{\boldsymbol{H_2}}$  and  $\boldsymbol{p}_{\boldsymbol{H_2O}}$  to satisfy both equations (122) and (123). (An alternate method of obtaining a solution of the simultaneous equations (122) and (123) is to plot each of these two equations as  $p_{H_2}$  vs  $p_{H_{20}}$  on the same graph and determine the point of intersection of the two curves. It may be found that in some cases more rapid convergence is obtained by this method.)

TABLE 1

Calculation of  $p_{H_2}$  and  $p_{H_20}$  from equations (122) and (123)

 $P_2 = 20 \text{ atmos}$   $T_2 = 3800 \,^{\circ}\text{K};$   $K_{p_{(H_20)}} = .3758 \text{ (atmos)} \frac{1}{2};$   $K_{p_{(H_2)}} = 1.107 \text{ (atmos)} \frac{1}{2};$   $K_{p_{(0_2)}} = 1.032 \text{ (atmos)} \frac{1}{2}.$ 

	Trial l	Trial 2	Trial 3	Trial 4
P <sub>H2</sub>	4. 35	3,00	3.43	3. 51
c <sub>ii</sub>	.0074634	.0156918	.012003	.0114640
C12	1.37607	1.47476	1.4362	1.42989
C <sub>13</sub>	-13.3412	-15.0826	-14.520	-14.4160
k <sub>ii</sub>	.0149268	.0313836	.024006	.022928
k <sub>IZ</sub>	. 232611	.302019	. 27463	. 270193
k <sub>is</sub>	-5,50442	- 3. 95870	-4, 4551	-4.54698
p <sub>H20</sub> (Eq. (122))	9. 23306	9. 30578	9. 3768	9. 3733
P <sub>HgO</sub> (Eq. (123))	12.9320	7.4068	9.0550	9. 3732
Δp <sub>H2O</sub>	+3.6989	-1.8990	3218	0001

Using equations (106) through (109), respectively, to calculate the values of the partial pressures of the remaining four constituents,

$$P_{02} = \left(\frac{K_{p_{(Hg0)}} p_{Hg0}}{p_{Hg}^2}\right)^2 = \left[\frac{(.3758)(9.3732)}{3.51}\right]^2 = 1.00710 \text{ atmos,}$$

$$p_{OH} = \frac{K_{P(H_2O)}^{I} p_{H_2O}}{p_{H_2}^{I}} = \frac{(.5984)(9.3732)}{(3.51)^{1/2}} = 2.99382 \text{ atmos},$$

$$p_{H} = K p_{(H_{2})} p_{H_{2}}^{\frac{1}{2}} = (1.107) (3.51)^{\frac{1}{2}} = 2.07396 \text{ atmos},$$

$$P_0 = \frac{K_{P(o_2)} K_{P(H_2O)} P_{H_2O}}{P_{H_2}} = \frac{(1.032)(.3758)(9.3732)}{3.51} = 1.03566 \text{ atmos}.$$

Rounding off to three decimal places all of the values of partial pressure obtained above, and adjusting the largest of these values ( $p_{H20}$ ) slightly in order to give a total ( $p_2$ ) of 20.000, the following values of partial pressure ( $p_i$ 's) are obtained:

$$\begin{array}{rcl}
 & P_{H_2} & = 3.510 \text{ atmos} \\
 & P_{H_20} & = 9.379 \\
 & P_{O_2} & = 1.007 \\
 & P_{O_1} & = 2.994 \\
 & P_{H_2} & = 2.074 \\
 & P_{O_2} & = 1.036 \\
 & P_{O_2} & = 20.000 \text{ atmos} \\
 & P_{O_2} & = 20.0000 \text{ atmos} \\
 & P_{O_2} & = 20.000 \text{ atmos} \\
 & P_{O_2} & = 20.0000 \text{ atm$$

Step 13. Total Number of Moles of the Final Mixture

Using equation (75),

$$n_2 = p_2 - \frac{\sum_{i}^{n_i} M_{w(i)}}{\sum_{j}^{p_j} M_{w(j)}}$$

$$= p_{2} \left[ \frac{n'_{H_{2}} M_{W(H_{2})} + n'_{0_{2}} M_{W(0_{2})}}{p_{H_{2}0} M_{W(H_{2}0)} + p_{H_{2}} M_{W(H_{2})} + p_{0_{2}} M_{W(0_{2})} + p_{0_{1}} M_{W(0H)} + p_{H} M_{W(H)} + p_{0} M_{W(0)}} \right]$$

or  $n_2 = 1.29676$ .

Step 14. Composition of the Final Mixture

Using equation (73) and values of p calculated in Step 12,

$$n_{H_2O} = \frac{p_{H_2O}}{p_e} n_z = \frac{(9.379)(1.29676)}{20.000} = .608116$$

$$n_{H_2} = \frac{p_{H_2}}{p_e} n_2 = \frac{(3.510)(1.29676)}{20.000} = .227581$$

$$n_{0g} = \frac{p_{0g}}{p_{s}} n_{g} = \frac{(1.007)(1.29676)}{20.000} = .065292$$

$$n_{OH} = \frac{p_{OH}}{p_{o}} n_{2} = \frac{(2.994)(1.29676)}{20.000} = .194125$$

$$n_H = \frac{p_H}{p_0} n_2 = \frac{(2.074)(1.29676)}{20.000} = .134474$$

$$n_0 = \frac{p_0}{p_2} n_2 = \frac{(1.036)(1.29676)}{20.000} = \frac{.067172}{1.29676}$$

#### Step 15. Properties of the Final Mixture

#### a. Mole Fractions

$$x_{H20} = \frac{n_{H20}}{n_2} = \frac{p_{H20}}{p_2} = \frac{9.379}{20.000} = .46895$$

$$x_{H2} = \frac{n_{H2}}{n_2} = \frac{p_{H2}}{n_2} = \frac{3.510}{20.000} = .17550$$

$$x_{O2} = \frac{n_{O2}}{n_2} = \frac{p_{O2}}{p_2} = \frac{1.007}{20.000} = .05035$$

$$x_{OH} = \frac{n_{OH}}{n_2} = \frac{p_{OH}}{p_2} = \frac{2.994}{20.000} = .14970$$

$$x_{H} = \frac{n_{H}}{n_2} = \frac{p_{H}}{p_2} = \frac{2.074}{20.000} = .10370$$

$$x_{O} = \frac{n_{O}}{n_2} = \frac{p_{O}}{p_2} = \frac{1.036}{20.000} = .05180$$

$$\sum_{j} x_{j} = \frac{1.00000}{1.00000}$$

#### b. Molecular Weight

$$M_{W_{2}} = \sum_{j} x_{j} M_{W(j)}$$

$$= x_{H_{2}0} M_{W(H_{2}0)} + x_{H_{2}} M_{W(H_{2})} + x_{O_{2}} M_{W(O_{2})}$$

$$+ x_{OH} M_{W(OH)} + x_{H} M_{W(H)} + x_{O} M_{W(O)}$$

$$= (.46895)(18.0160) + (.17550)(2.0160)$$

$$+ (.05035)(32.0000) + (.14970)(17.0080)$$

$$+ (.10370)(1.0080) + (.05180)(16.0000),$$

or

$$M_{W_2} = 13.8930.$$

#### c. Molal Heat Capacity\*

<sup>\*</sup> The molal heat capacities for H<sub>2</sub>O, H<sub>2</sub>, O<sub>2</sub> and OH were obtained from Figure 16 through 19, respectively; and the values for O and H were calculated from enthalpy data given in Ref. 7, page 22.

$$C_{p_2} = \sum_{j} x_j C_{p_{(j)}}$$

$$= x_{H_2O} C_{p_{(H_2O)}} + x_{H_2} C_{p_{(H_2)}} + x_{O_2} C_{p_{(O_2)}} + x_{OH} C_{p_{(OH)}}$$

$$+ x_H C_{p_{(H)}} + x_O C_{p_{(O)}}$$

or

$$C_{P_2} = 10.7339 \text{ cal/(gm mole)(°K)}.$$

d. Specific Heat Ratio

$$\chi_2 = \frac{C_{p_2}}{C_{p_2} - R} = \frac{10.7339}{10.7339 - 1.98719} = 1.2272.$$

e. Heat of Formation  $\Delta H'_{f(i)}$ , of Each Product at the Initial Temperature\*

By definition of  $\Delta H'_{(j)}$ , in the same manner as in Step 8 for the initial mixture, the heat of formation of an element (in its standard state of aggregation) at  $t_1$  is equal to zero. Therefore,

$$\Delta H'_{f(H_2)} = 0.0,$$

$$\Delta H^{i}_{f(o_2)} = 0.0.$$

From the literature, at  $t_{I} = 18^{\circ}C$ ,

$$\Delta H_{f_L (H_20)} = +57,801** cal/gm mole,$$

but, since  $t_{L} = t_{1}$  in this case,

$$\Delta H_{f(H_2O)}^{i} = \Delta H_{fL(H_2O)}$$
 = +57,801 cal/gm mole.

<sup>\*</sup> See footnote in Step 8, par. e. \*\* Ref. 2, page 255.

From the literature, at  $t_{1} = 25^{\circ}C$ ,

 $\Delta H_{f_L \text{ (OH)}} = -9,310* \text{ cal/gm mole,}$ 

 $\Delta H_{f_L(H)} = -51,900* cal/gm mole,$ 

 $\Delta H_{f_1}$  (o) = -59, 100\* cal/gm mole,

but, since  $t_{L} \neq t_{1}$  in this case, Equation (81) is used to calculate values of  $\Delta H_{f(i)}'$  for OH, H and O\*\*:

$$\Delta H'_{f(j)} = \Delta H_{f(j)} - H'_{j} + \sum H'_{e(j)}$$
 (81)

Values of  $H'_{j}$  and  $\sum H'_{e(j)}$  obtained from the literature for OH, H and O are indicated in Table 2.

From the above data and equation (81),

$$\Delta H'_{f(OH)} = \Delta H_{fL(OH)} - H'_{(OH)} + \sum H'_{e(OH)}$$
  
= -9,310 + 47.9 - 48.5 = -9310.6 cal/gm mole;

$$\Delta H_{f(H)}^{i} = \Delta H_{f_{L}(H)} - H_{(H)}^{i} + \sum H_{e(H)}^{i}$$
  
= -51,900+33.2 - 24.0 = -51,890.8 cal/gm mole;

$$\Delta H_{f(0)}' = \Delta H_{f(0)} - H_{(0)}' + \sum H_{e(0)}'$$
  
= -59, 100+33.2 - 24.5 = -59, 091.3 cal/gm mole.

<sup>\*</sup> Ref. 7, page 26.

<sup>\*\*</sup> The difference between  $\Delta H'_{f(j)}$  and  $\Delta H_{f_{L(j)}}$  for this case  $(t_L - t_1 = 7^{\circ}C)$  is negligible (<0.1%). The calculation of  $\Delta H'_{f(j)}$  according to equation (81) has been included for purpose of illustration.

Literature Values of Enthalpies, cal/gm mole of "j"

TABLE 2

Component "j"	H'; (cal/gm mole)	$\Sigma_{H_{e}(j)}$ (cal/gm mole of "j")	Reference
ОН	-47. 9	-48.5	Calculated (*) from
н	-33.2	-24.0	data in Ref. 7,
0	-33. 2	-24.5	page 22 and Ref. 2,
			page 216.

$$H'_{(OH)} = \int_{T_L}^{T_I} \left(\frac{dH}{dT}\right)_{(OH)} dT = \int_{T_L}^{T_I} C_{P_{(OH)}} dT \cong \overline{C}_{P_{(OH)}} \left(T_I - T_L\right), \qquad (124)$$

where  $\overline{C}_{p_{(0H)}}$  = mean  $C_p$  between  $T_L$  (298°K) and  $T_1$  (291°K) = 6.8421  $\frac{\text{cal}}{(\text{gm mole})(^{\circ}\text{K})}$ ; therefore,  $H'_{(0H)}$  = (6.8421)(291-298) = -47.9 cal/gm mole.

The formation reaction for one mole of OH is  $1/2 H_2 + 1/2 O_2 \longrightarrow$  OH; then,

$$\Sigma H'_{e_{(OH)}} = n_{H_2} \overline{C}_{P_{(H_2)}} (T_I - T_L) + n_{O_2} \overline{C}_{P_{(O_2)}} (T_I - T_L)$$
,

where  $\overline{C}p_{(H_2)}$  = mean  $C_p$  between  $T_L$  (298°K) and  $T_1$  (291°K) = 6.86 cal/(gm mole)(°K), and  $\overline{C}p_{(0_2)}$  = mean  $C_p$  between  $T_L$  (298°K) and  $T_1$  (291°K) = 7.00 cal/(gm mole)(°K);

$$\Sigma H'_{e (OH)}$$
 = (1/2)(6.86)(291-298)+1/2(7.00)(291-298) = -48.5  $\frac{\text{cal}}{\text{gm mole of OH}}$ .

The values of the heats of formation,  $\Delta H'_{f(j)}$ , calculated above, are summarized for convenience in Table 3.

<sup>\*</sup> The calculation of  $H'_j$  and  $\sum H'_{e(j)}$  from data obtained in References 2 and 7 is shown below for OH.

TABLE 3

#### Heats of Formation of Products, cal/gm mole

ΔH' <sub>f(H2</sub> )	ΔH' <sub>f(O2</sub> )	ΔH' <sub>f (H20)</sub>	∆ H' <sub>f (он)</sub>	Δ H' <sub>f (H)</sub>	ΔH' <sub>f(o)</sub>
0.0	0.0	+57,801	-9310.6	-51,890,8	-59,091.3

Step 16. Heat of Reaction, Energy Release per Mole of Initial Mixture

Using data from Step 14 and Step 15 (par. e),

$$\begin{split} \sum_{j} n_{j} \triangle H_{f(j)}^{i} &= n_{\text{H2O}} \triangle H_{f(\text{H2O})}^{i} + n_{\text{H2}} \triangle H_{f(\text{H2})}^{i} + n_{\text{O2}} \triangle H_{f(\text{O2})}^{i} + n_{\text{OH}} \triangle H_{f(\text{OH})}^{i} + n_{\text{H}} \triangle H_{f(\text{H})}^{i} + n_{\text{O}} \triangle H_{f(\text{O})}^{i} \\ &= (.608116)(+57,801) + (.227581)(0.0) + (.065292)(0.0) \\ &+ (.194125)(-9310.6) + (.134474)(-51,890.8) + (.067172) \\ &- (-59,091.3) \\ &= +22,395.0 \text{ calories}; \\ \sum_{i} n_{i}^{i} \triangle H_{f(i)}^{i} &= n_{\text{H2}}^{i} \triangle H_{f(\text{H2})}^{i} + n_{\text{O2}}^{i} \triangle H_{f(\text{O2})}^{i} \\ &= (1) (0.0) + (1/2) (0.0) = 0; \\ n_{1} &= n_{\text{H2}}^{i} + n_{\text{O2}}^{i} = 1 + \frac{1}{2} = 1.5 \text{ moles} \; , \end{split}$$

then, by equation (83),

$$\Delta H_{R}$$
 =  $\frac{\sum_{j} n_{j} \Delta H'_{f(j)} - \sum_{i} n'_{i} \Delta H'_{f(i)}}{n_{1}} = \frac{+22,395,0-0.0}{1.5}$ ,

or  $\Delta H_R = +14,930.0$  cal/gm mole of initial mixture.

#### Step 17. Energy Release Function

Using equation (84) and data from Steps 7, 8 (par. d) and 16,

$$B = \frac{\Delta H_R}{\chi_{1} RT_{1}} = \frac{+14,930.0}{(1.4026)(1.98719)(291)} = 18.407.$$

#### Step 18. Gamma Ratio

From Figure 5 and data from Steps 8 (par. d.) and 15 (par. d.) or by direct calculation,

$$\frac{\mathbf{\chi}_2}{\mathbf{\chi}} = \frac{1.2272}{1.4026} = .87495.$$

#### Step 19. Molecular Weight Ratio

Using data from Steps 8 (par. b ) and 15 (par. b ),

$$\frac{\mathbf{M_{w_2}}}{\mathbf{M_{w_1}}} = \frac{13.8930}{12.0107} = 1.15672.$$

### Step 20. The Dimensionless Parameter E

From Figure 1 and data from Steps 15 (par. d) and 17, or using equation (47).

$$E = (\chi_2^2 - 1)B = [(1.2272)^2 - 1] [18.407] = 9.31431$$
.

#### Step 21. Gamma Function

From Figure 2 and data from Steps 8 (par. d) and 15 (par. d), or using equation (48),

$$\Gamma = \frac{\chi_2^2 - \chi_1}{\chi_1(\chi_1 - 1)} = \frac{(1.2272)^2 - 1.4026}{1.4026(1.4026-1)} = 0.183146.$$

#### Step 22. Pseudo-Mach Number (Squared)

From Figure 3 and data from Steps 20 and 21, or using equation (45),

$$\overline{M}_1^2 = 2E + 2\Gamma = (2)(9.31431) + 2(0.183146) = 18.99491$$
.

#### Step 23. Mach Number (Squared)

From Figure 4, and values of  $\overline{M}_1^2$  and  $\frac{\sqrt{02}}{\sqrt{01}}$  from Steps 18 and 22,  $M_1^2 = 19.0$  (Same as  $\overline{M}_1^2$  within the accuracy of reading Figure 4).

### Step 24. The Dimensionless Parameter $\sqrt[7]{M_1^2}$

From Figure 9 and data from Steps 8 (par. d) and 23, or by direct calculation.

$$\sqrt[8]{M_1^2} = (1.4026)(19.0) = 26.65$$
.

#### Step 25. Temperature - Molecular Weight Ratio

From Figure 8 and data from Steps 15 (par. d) and 24, or using equation (85),

$$\frac{T_2 M_{W_1}}{T_1 M_{W_2}} = \frac{\sqrt[8]{2}}{(\sqrt[8]{2}+1)^2} \left(\sqrt[8]{1} M_1^2 + \frac{1}{\sqrt[8]{1} M_1^2} + 2\right)$$

$$= \frac{1.2272}{(1.2272+1)^2} (26.65 + \frac{1}{26.65} + 2)$$

$$= 7.0972.$$

#### Step 26. Final Temperature

Using data from Steps 7, 19 and 25,

$$T_2 = \frac{T_2 M_{w_1}}{T_1 M_{w_2}} \cdot \frac{M_{w_2}}{M_{w_1}} \cdot T_1$$
  
= (7.0972)(1.15672)(291)  
= 2389°K.

Since this calculated value of  $T_2$  did not check the assumed value (Step 10), other values of  $T_2$  were assumed and Steps 10 through 26 were repeated until sufficient data were obtained to permit graphical determination of the correct value of  $T_2$  corresponding to the assumed pressure.

A plot of  $\Delta$   $p_{\rm H_2O}$  vs  $p_{\rm H_2}$  (assumed) for each assumed  $T_2$  is shown in Figure 21, and a plot of  $T_2$  (calculated) vs  $T_2$  (assumed) is shown in Figure 22 (For Trials 1 through 29, i.e., for  $p_2$  (assumed) = 20.00 atmos). Table 4 lists principal intermediate values calculated in these trials.

The correct value of  $T_2$  corresponding to the first assumed  $p_2$  of 20.00 atmos is shown in Figure 22 to be 3565°K (the point of intersection of the curve with the 45° line). This value of  $T_2$  is also indicated in Table 4 (Trials 25-29) as  $T_2$  (adjusted).

It was considered that the value of  $T_{2(adjusted)}$  obtained in Trial 29 as described above was so close to the value of  $T_{2(assumed)}$  that further calculations were not warranted. This was considered to be justifiable since the thermodynamic data ( $K_p$  and  $C_p$ ) available were not of sufficient accuracy to give data different from that obtained in the last trial (Trial 29). Hence, all intermediate values calculated in the last trial (Trial 29), based on the assumed  $T_2$  for that trial, for all practical purposes can be considered to be accurate values corresponding to the adjusted  $T_2$ .

#### Step 27. Pressure Ratio

Using the values of  $\tilde{b}_2$  and  $\tilde{b}_1 M_1^2$  from Trial 29, then from Figure 6 or using equation (87),

$$\frac{p_2}{p_1} = \frac{1}{Y_2 + 1} \left( Y_1 M_1^2 + 1 \right)$$

$$= \frac{1}{1.20714 + 1} (37.3793 + 1)$$

$$= 17.3887.$$

#### Step 28. Final Pressure

Using data from Steps 6 and 27,

$$p_2 = \frac{p_2}{p_1} \cdot p_1 = (17.3887)(1) = 17.3887 \text{ atmos.}$$

Since this calculated value of  $p_2$  did not check the assumed value (Step 9), other values of  $p_2$  were assumed and Steps 9 through 28 were repeated until sufficient data were obtained to permit graphical determination of the correct value of  $p_2$ .

A plot of  $\Delta p_{\text{H}_2\text{O}}$  vs  $p_{\text{H}_2(\text{assumed})}$  for each assumed  $T_2$  is shown in Figures 21, 23, 25, 26 for assumed  $p_2$ 's of 20, 17, 17.17 and 17.20, respectively. A plot of  $T_2(\text{calculated})$  vs  $T_2(\text{assumed})$  is shown in Figures 22, 24 and 27 for assumed  $p_2$ 's of 20, 17 and 17.20, respectively\*.

<sup>\*</sup> For an assumed p<sub>2</sub> of 17.17, no plot of T<sub>2</sub>(calculated) vs T<sub>2</sub>(assumed) was necessary since the calculated T<sub>2</sub> for the first trial was so close to the assumed T<sub>2</sub> that additional trials were not warranted.

Figure 28 shows a plot of p<sub>2</sub>(calculated) vs p<sub>2</sub>(assumed) Table 4 lists principal intermediate values calculated in Trials 1 through 53.

The correct value of the final detonation pressure, p<sub>2</sub>, is shown in Figure 28 to be 17.22 atmos (the point of intersection of the curve with the 45° line). This value of p<sub>2</sub> is also indicated as p<sub>2</sub>(adjusted) in Table 4.

It was considered that the value of  $p_{2(adjusted)}$  obtained in Trial 53 as described above was so close to the value of  $p_{2(assumed)}$  that further calculations were not warranted. Hence, all intermediate values calculated in the last trial (Trial 53), based on the assumed  $p_{2}$  for that trial, can be considered to be accurate values corresponding to the adjusted  $p_{2}$ .

## Step 29. Detonative Mach Number

Using the value of M<sub>1</sub><sup>2</sup> from Trial 53,

$$M_1 = \sqrt{M_1^2} = \sqrt{26.40} = 5.13809$$
.

# Step 30. Detonation Velocity

Using equation (89) and data from Steps 7, 8 (pars. b and d) and 29,

$$D = M_1 \sqrt{\frac{J \gamma_1 R T_1}{M_{W_1}}}$$

$$= 5.13809 \sqrt{\frac{(4.185 \times 10^7)(1.4026)(1.98719)(291)}{12.0107}}$$

$$= 2.73149 \times 10^5$$
 cm/sec.

This detonation velocity expressed in more common units is

or

#### Step 31. Calculation of Additional Data

### a. Density Ratio

Using the values of  $\frac{P_2}{P_1}$  and  $\frac{T_2M_{W_1}}{T_1M_{W_2}}$  from Trial 53,

then from Figure 10 or using equation (15),

$$\frac{P_2}{P_1} = \frac{P_2}{P_1} / \frac{T_2 M_{W_1}}{T_1 M_{W_2}} = \frac{17.2269}{9.67764} = 1.78007.$$

### b. Initial Density

Using equation (90) and the values of  $P_1$ ,  $T_1$  and  $M_{W_1}$  from Steps 6, 7 and 8 (par. b ) of this sample calculation,

$$\rho_{1} = \frac{M_{W_{1}}P_{1}}{JRT_{1}} = \frac{(12.0107)(1.0133\times10^{6})}{(4.185\times10^{7})(1.98719)(291)} = 5.02895\times10^{-4} gm/cm_{3}^{3}$$

or initial density is 0.5029 gm/liter.

## c. Final Density

Using the above values of  $\frac{\rho_2}{\rho_1}$  and  $\rho_1$ ,

$$P_2 = \frac{P_2}{P_1} \cdot P_1 = (1.78007)(5.02895 \times 10^{-4})$$
  
= 8.95188 × 10<sup>-4</sup> gm/cm<sup>3</sup>,

or final density is 0.8952 gm/liter.

### a. Density Ratio

Using the values of  $\frac{P_2}{P_1}$  and  $\frac{T_2M_{W_1}}{T_1M_{W_2}}$  from Trial 53,

then from Figure 10 or using equation (15),

$$\frac{\rho_2}{\rho_1} = \frac{P_2}{P_1} / \frac{T_2 M_{W_1}}{T_1 M_{W_2}} = \frac{17.2269}{9.67764} = 1.78007.$$

## b. Initial Density

Using equation (90) and the values of  $P_1$ ,  $T_1$  and  $M_{w_1}$  from Steps 6, 7 and 8 (par. b ) of this sample calculation,

$$\rho_{l} = \frac{M_{W_{1}}P_{1}}{JRT_{1}} = \frac{(12.0107)(1.0133\times10^{6})}{(4.185\times10^{7})(1.98719)(291)} = 5.02895\times10^{-4} \text{gm/cm}_{3}^{3}$$

or initial density is 0.5029 gm/liter.

### c. Final Density

Using the above values of  $\frac{\rho_2}{\rho_i}$  and  $\rho_i$ ,

or final density is 0,8952 gm/liter.

2

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β E¢ (20)

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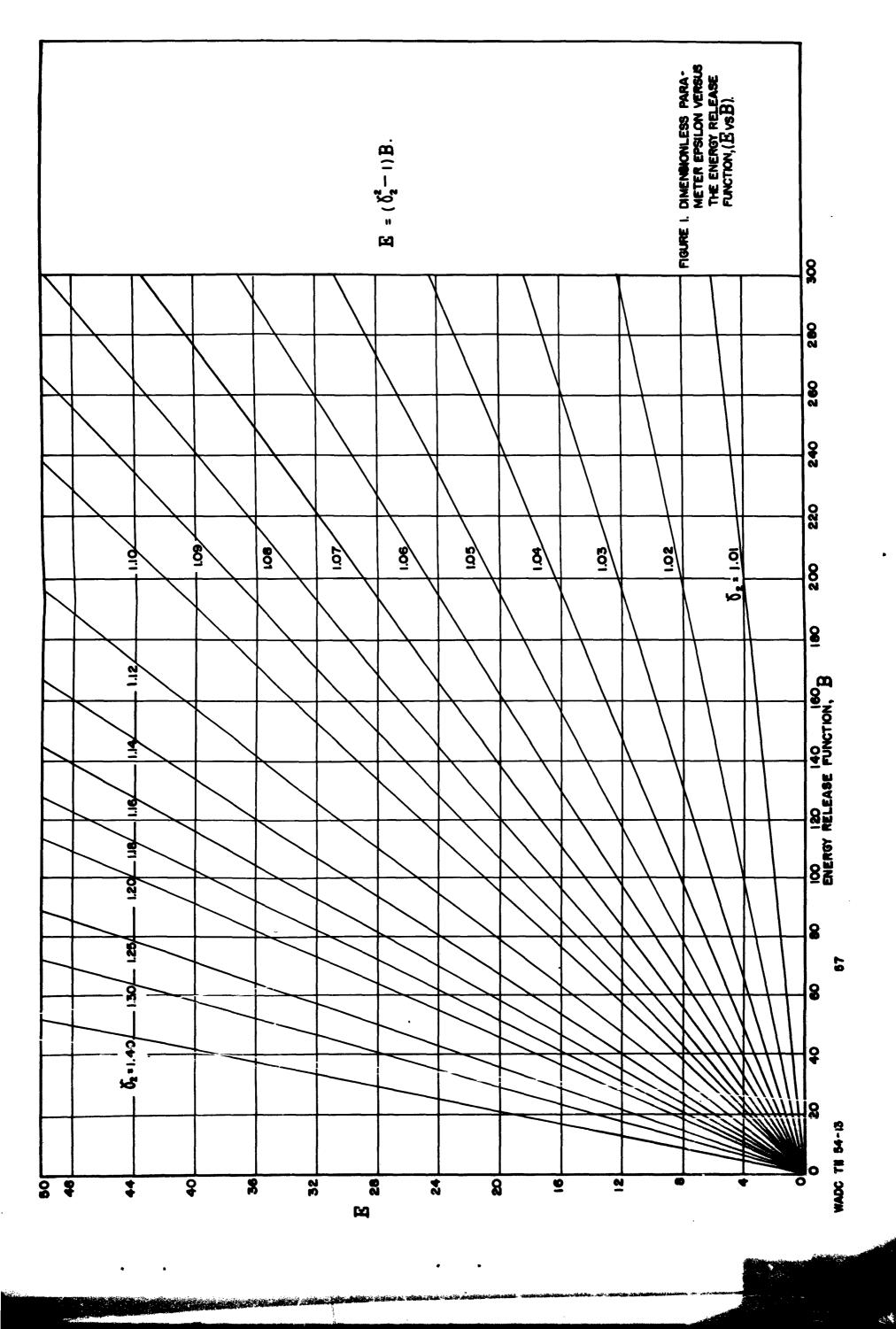
APPENDIX IV

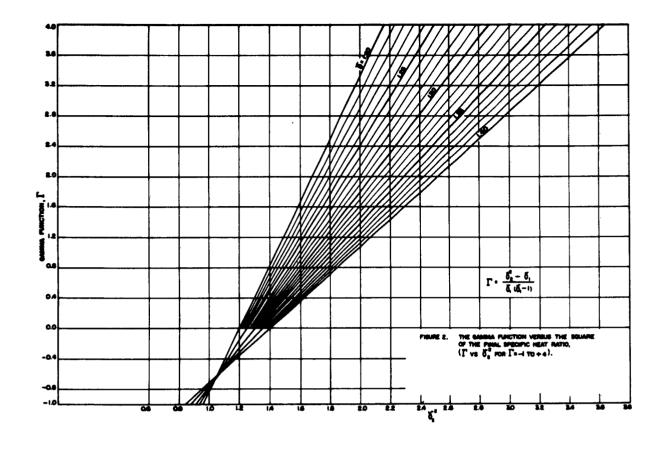
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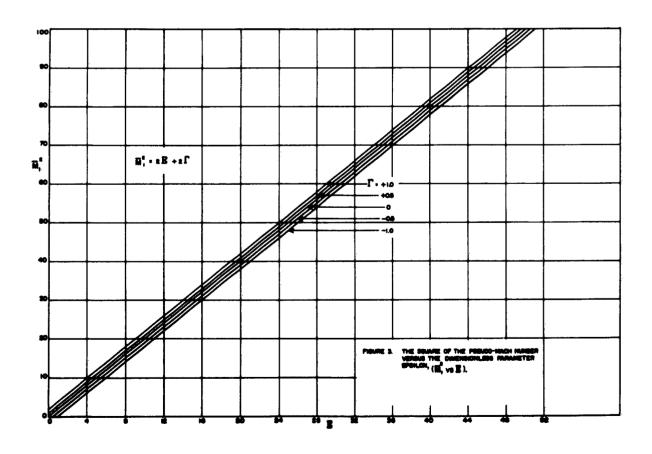
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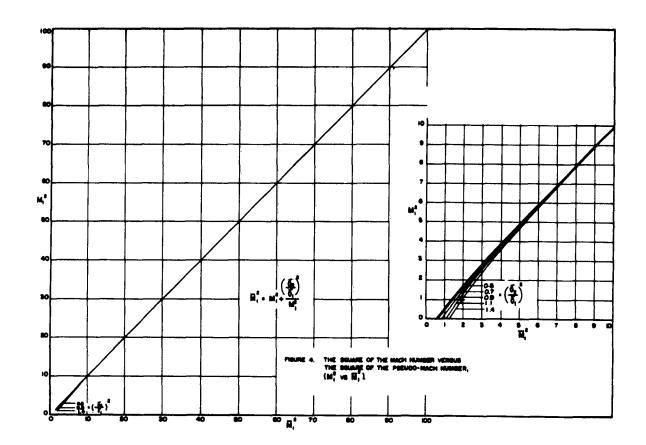
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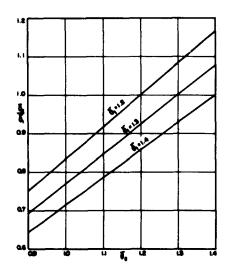
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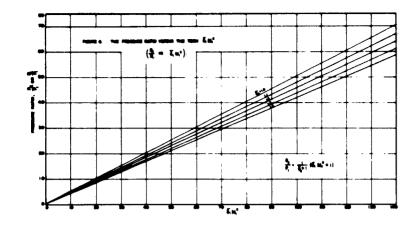


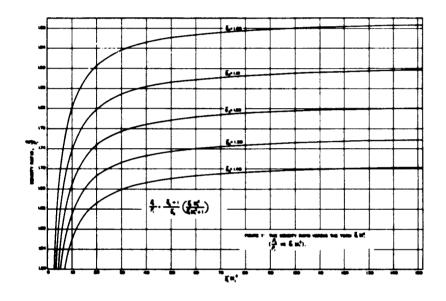


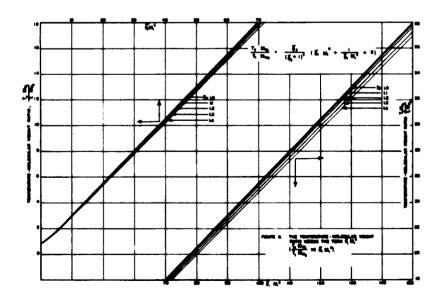




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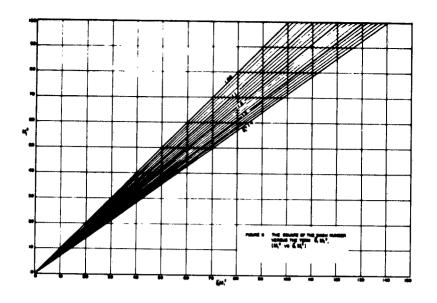


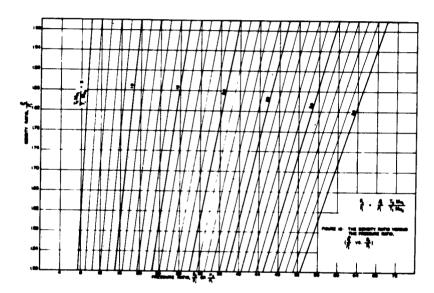


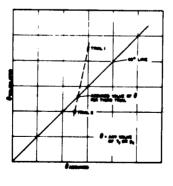


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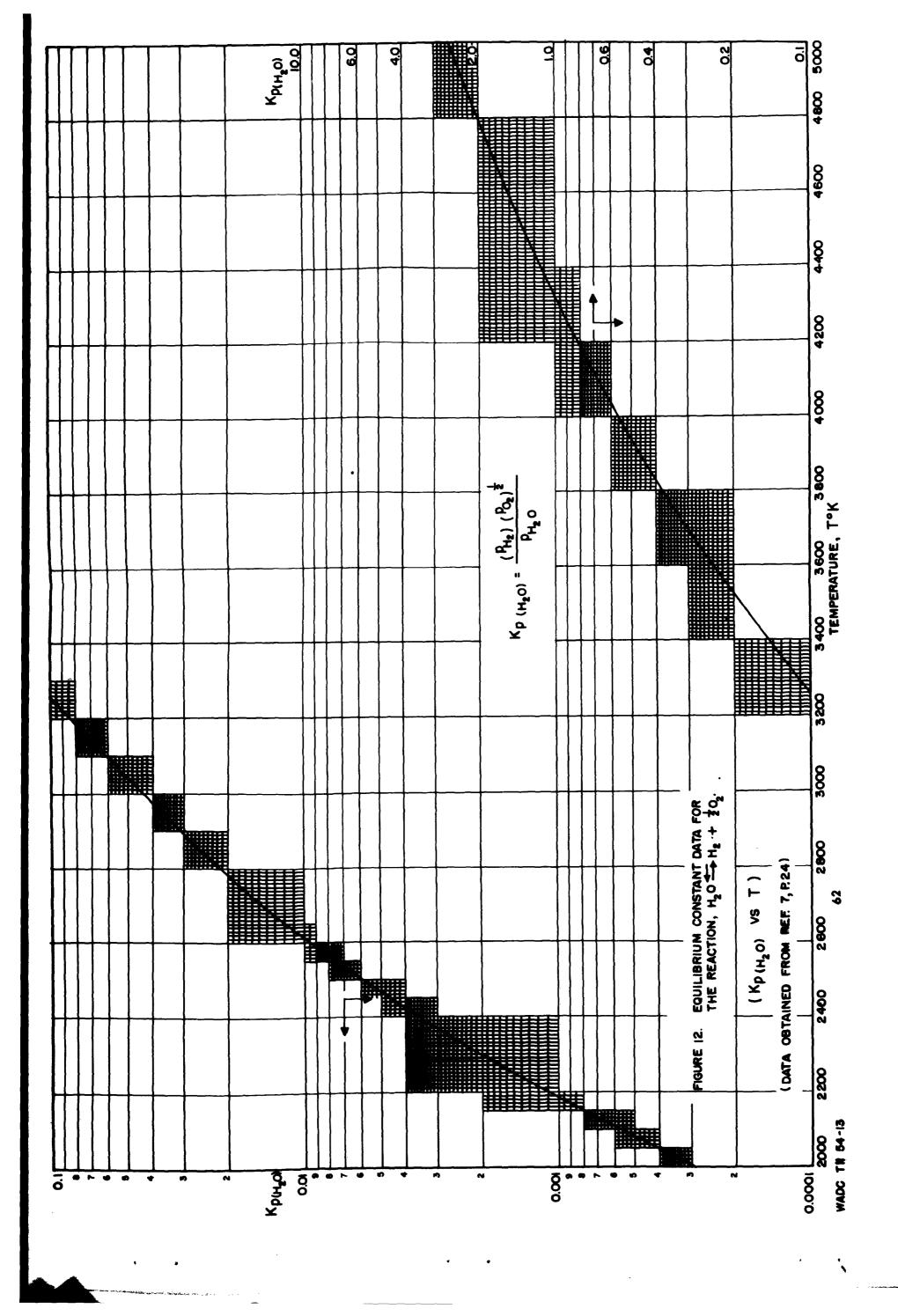


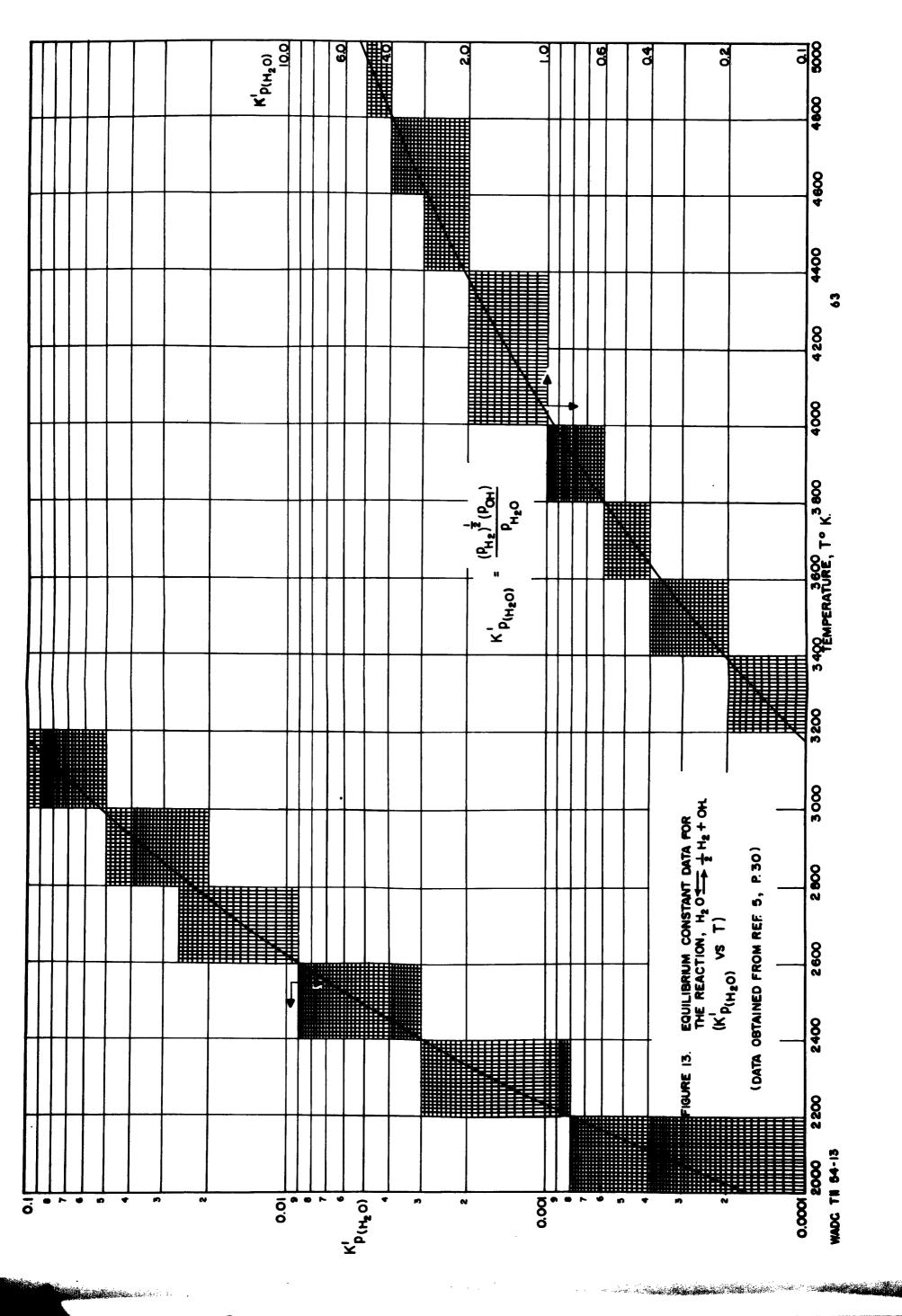


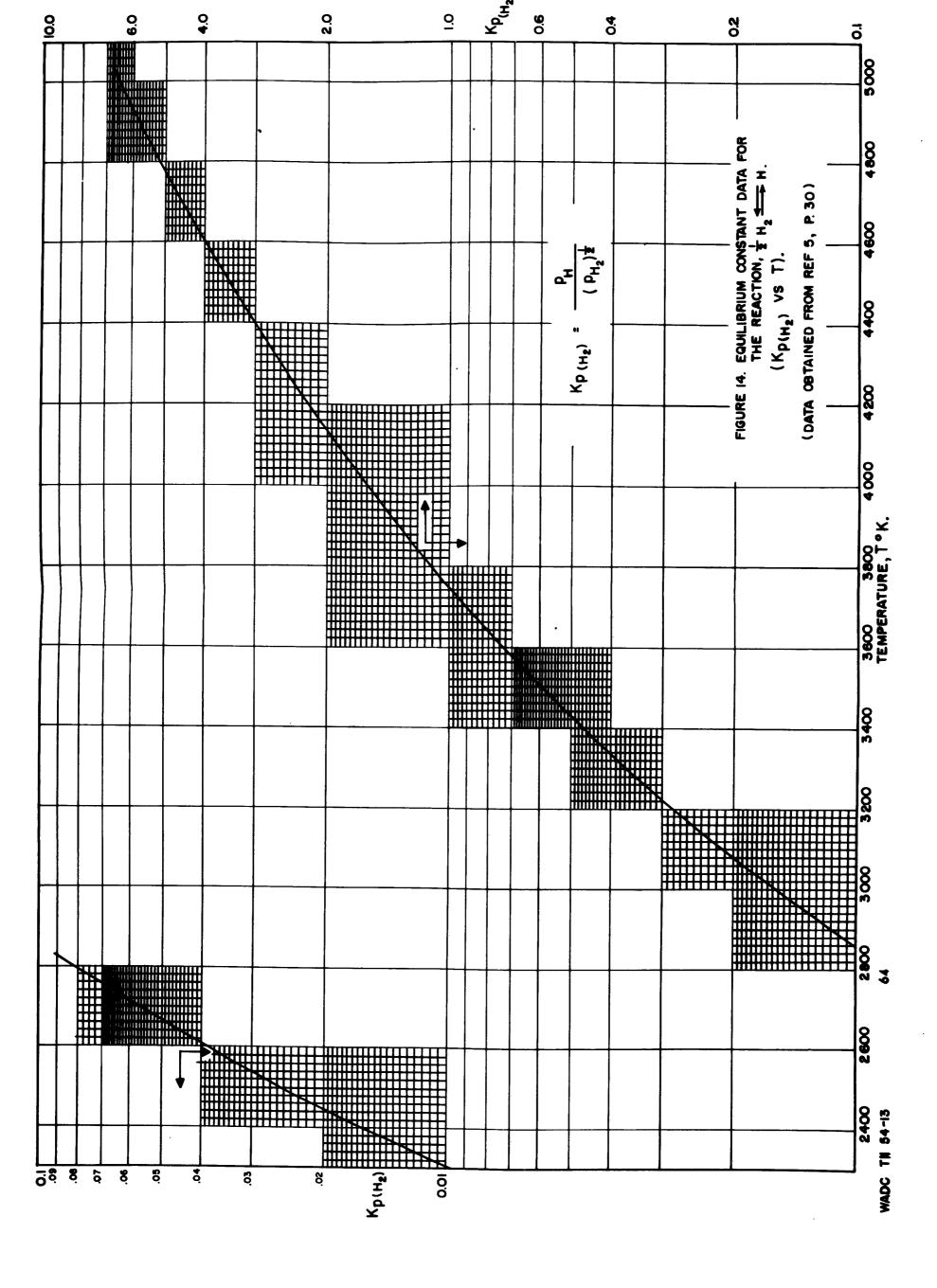


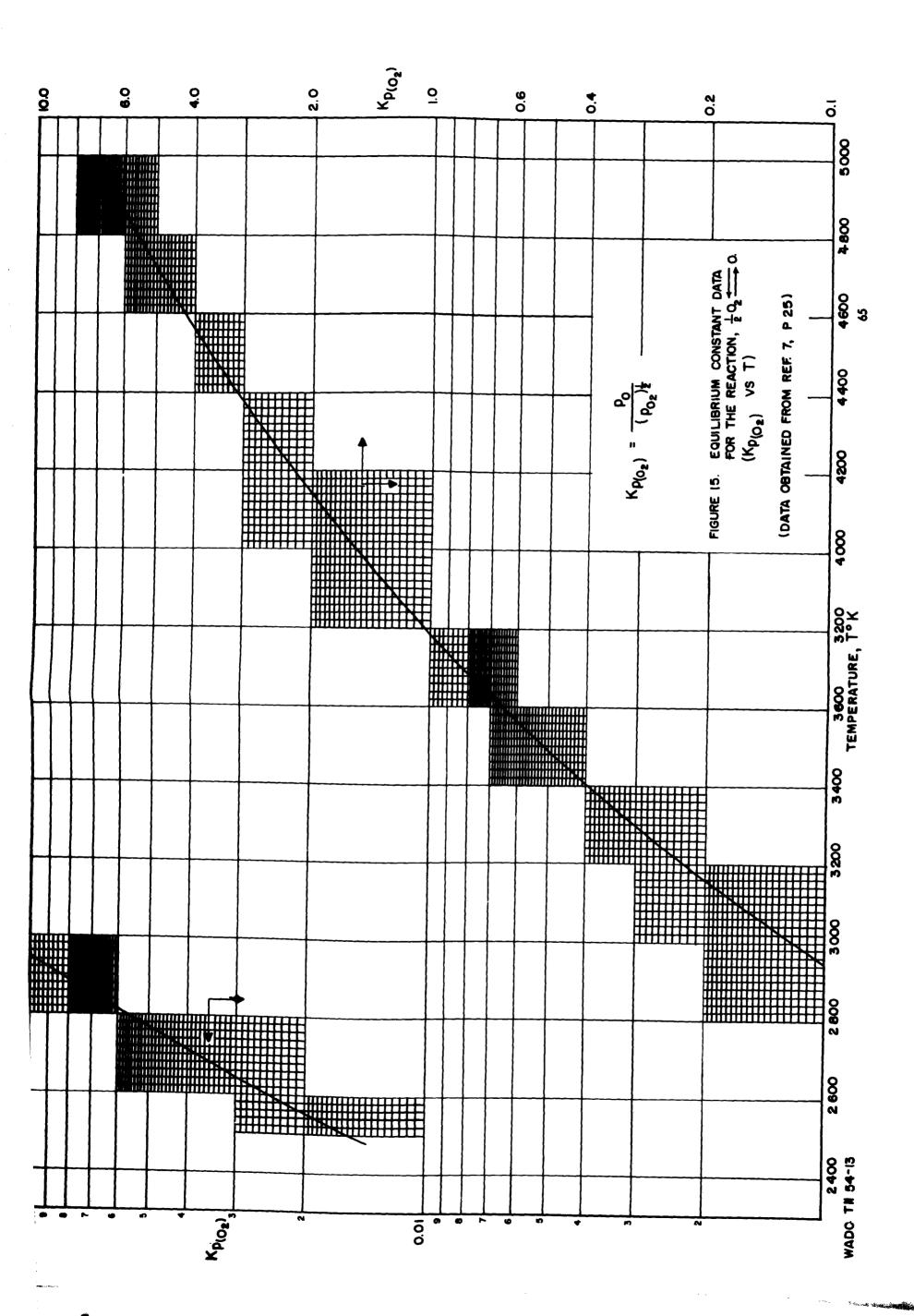
Posite II. Threat pur par many compresses or their our-parties obtained.

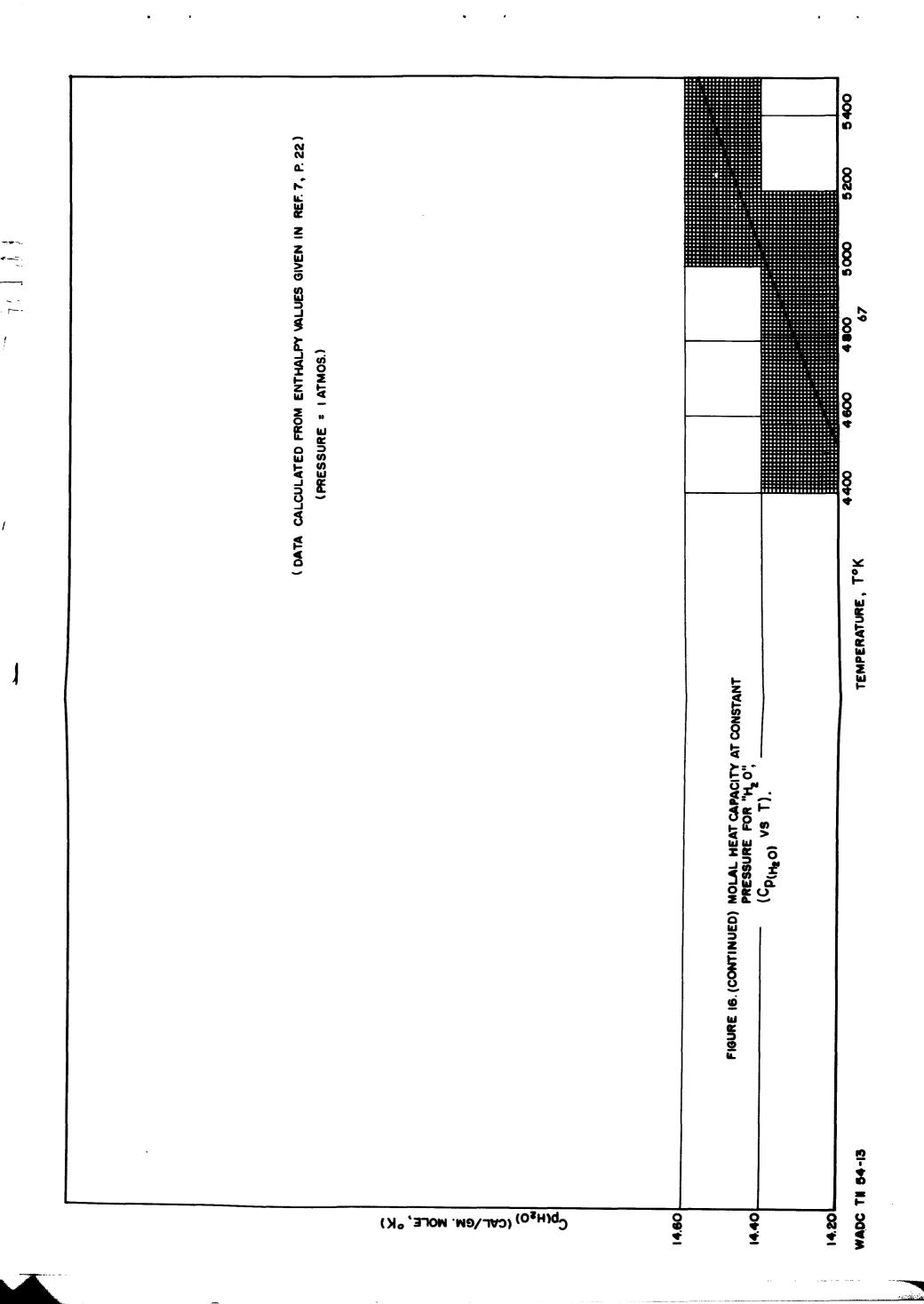
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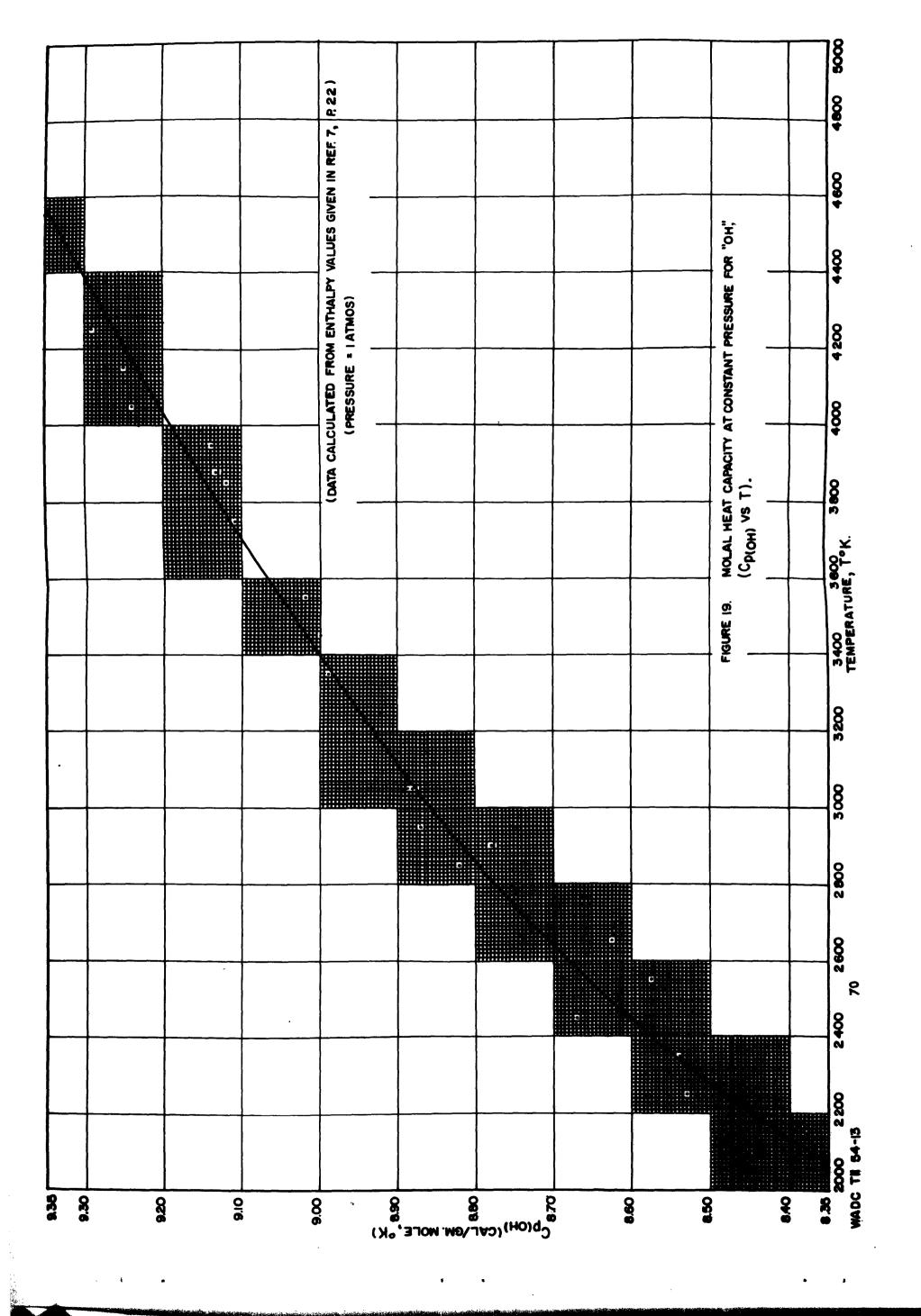




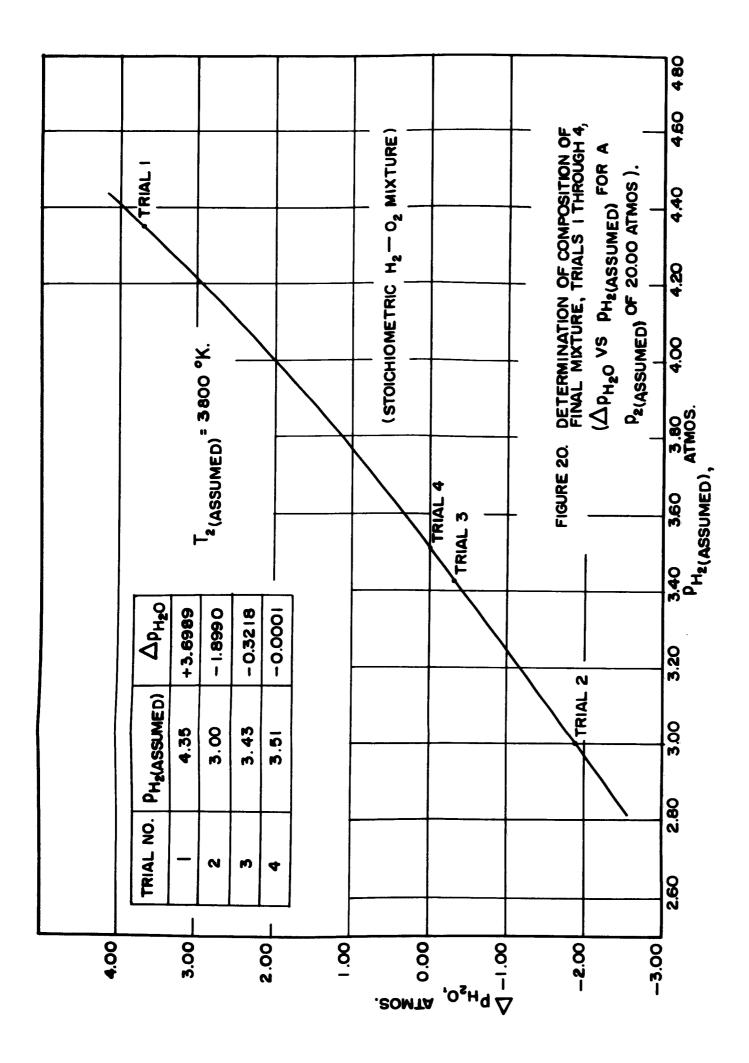




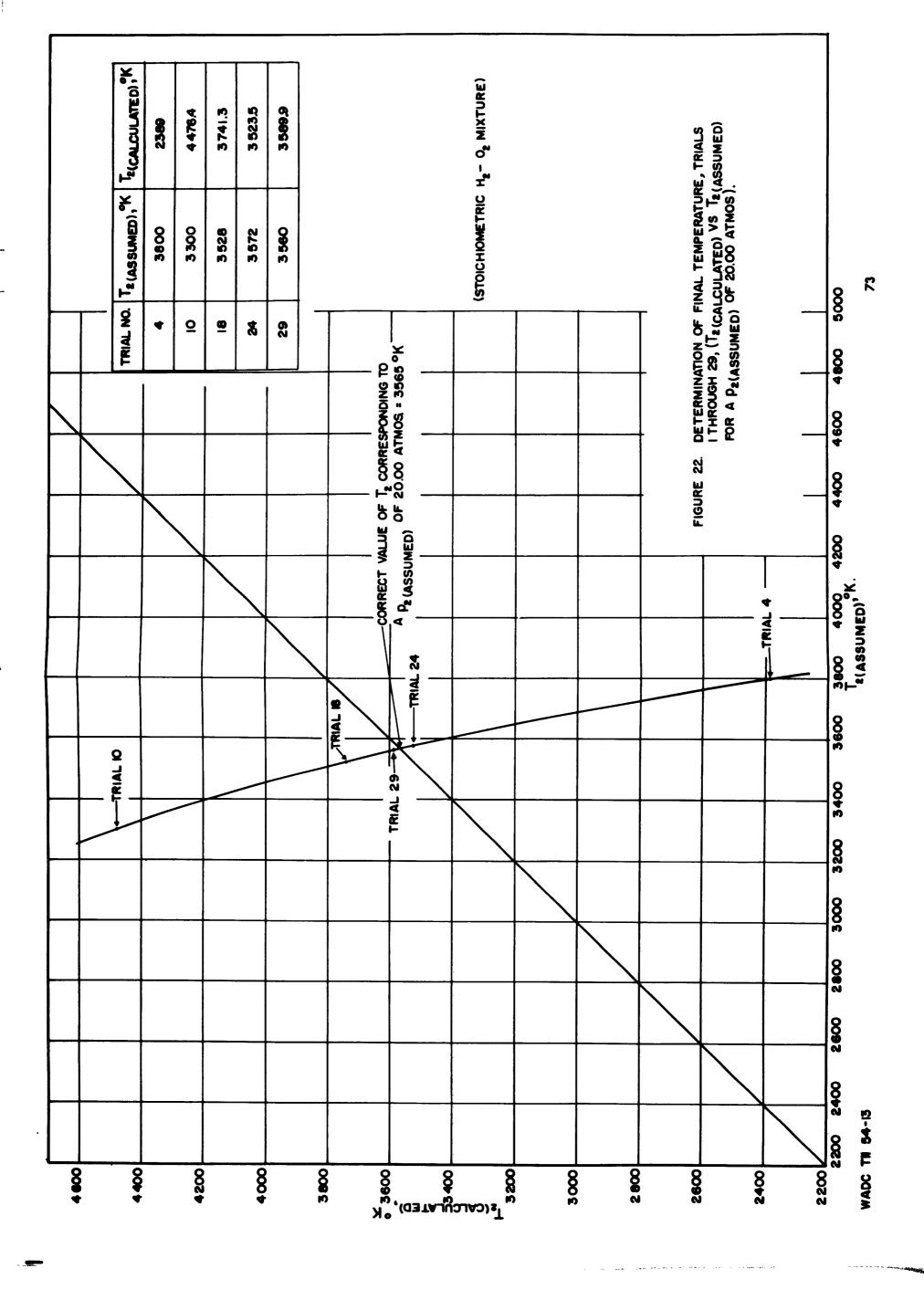








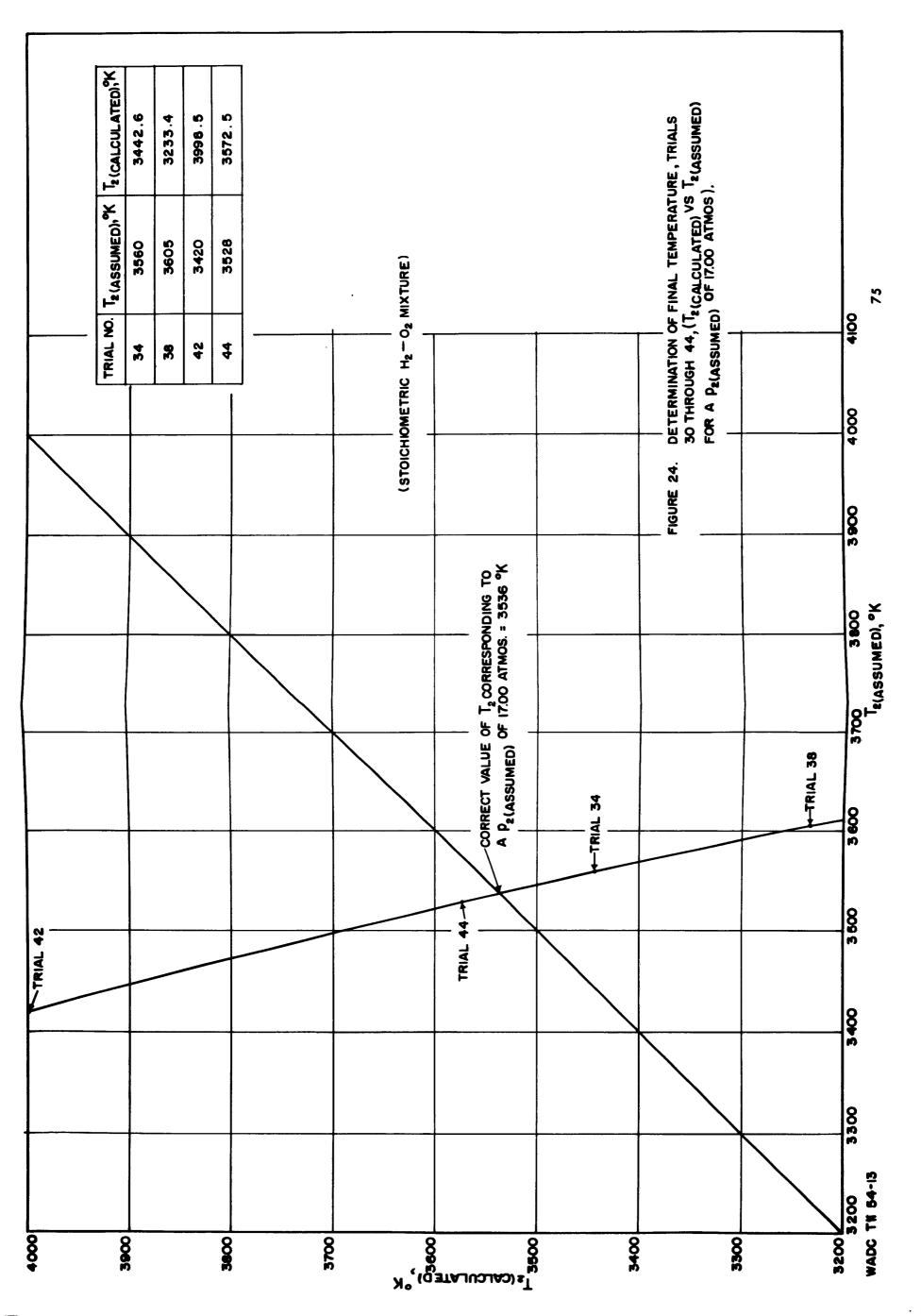
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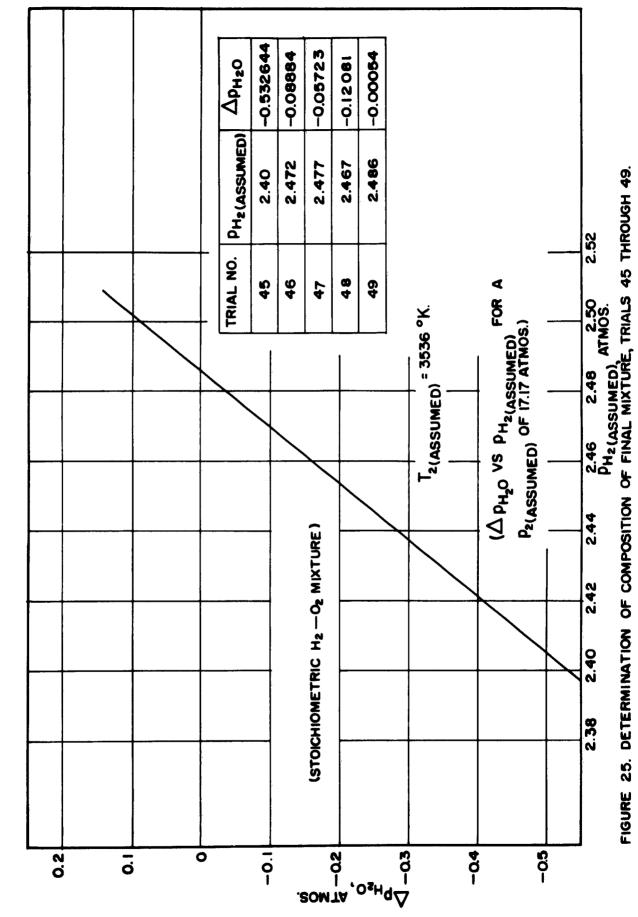


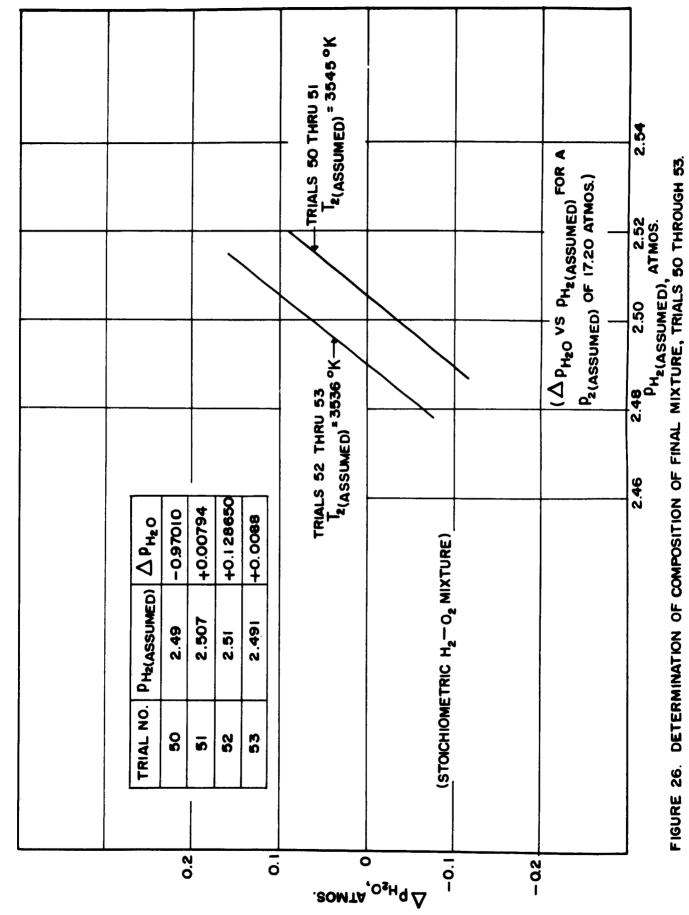
<del></del>
TRIALS 43 THRU 44  TE(ASSUMED) = 35  TRIALS 35  TRIALS 35  TRIALS 35  FIGURE 23.
TRIAL NO.  TRIALS 30 THRU 34  TRIALS 30 THRU 34  TRIALS 35 THRU 34  TRIALS 35 THRU 38  TRIALS 35 THRU 44  (\$\triangle \text{DPH2} \text{OVS} \text{PPRUMED})  PE(ASSUMED) OF IT.OO ATMOS
TRIAL NO. 30 31 32 33 34 35 36 36 37 38 39 40 41 42 43 44 42 43 44 44 42 43 44 44 44 44 44 44 44 44 44 44 44 44
PHz(ASSUMED) 2.50 2.80 2.649 2.537 2.526 2.66 2.66 2.66 2.66 2.19 2.19 2.15 2.14 2.15 2.14 2.15 2.40 2.44 2.15 4, FOR A

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